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Slide of the Seminar

How quickly does turbulence die out ?

Dr. Gregory Bewley

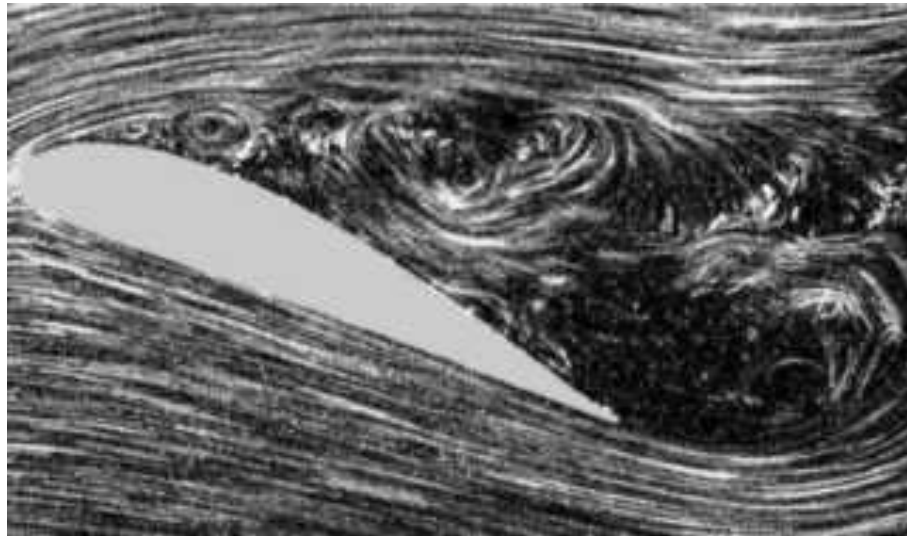
***ERC Advanced Grant (N. 339032) "NewTURB"
(P.I. Prof. Luca Biferale)***

Università degli Studi di Roma Tor Vergata
C.F.n. 80213750583 – Partita IVA n. 02133971008 - Via della Ricerca Scientifica, 1 – 00133 ROMA

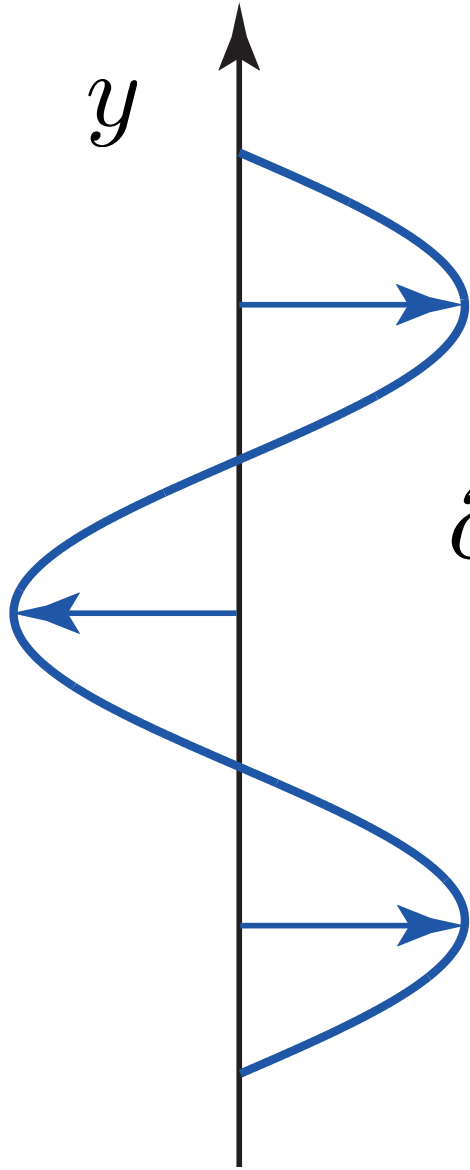
How quickly does turbulence die out?

Gregory P. Bewley
Michael Sinhuber
Eberhard Bodenschatz

Max Planck Institute for Dynamics and
Self-Organization
Göttingen, Germany



http://en.wikipedia.org/wiki/File:Flow_separation.jpg



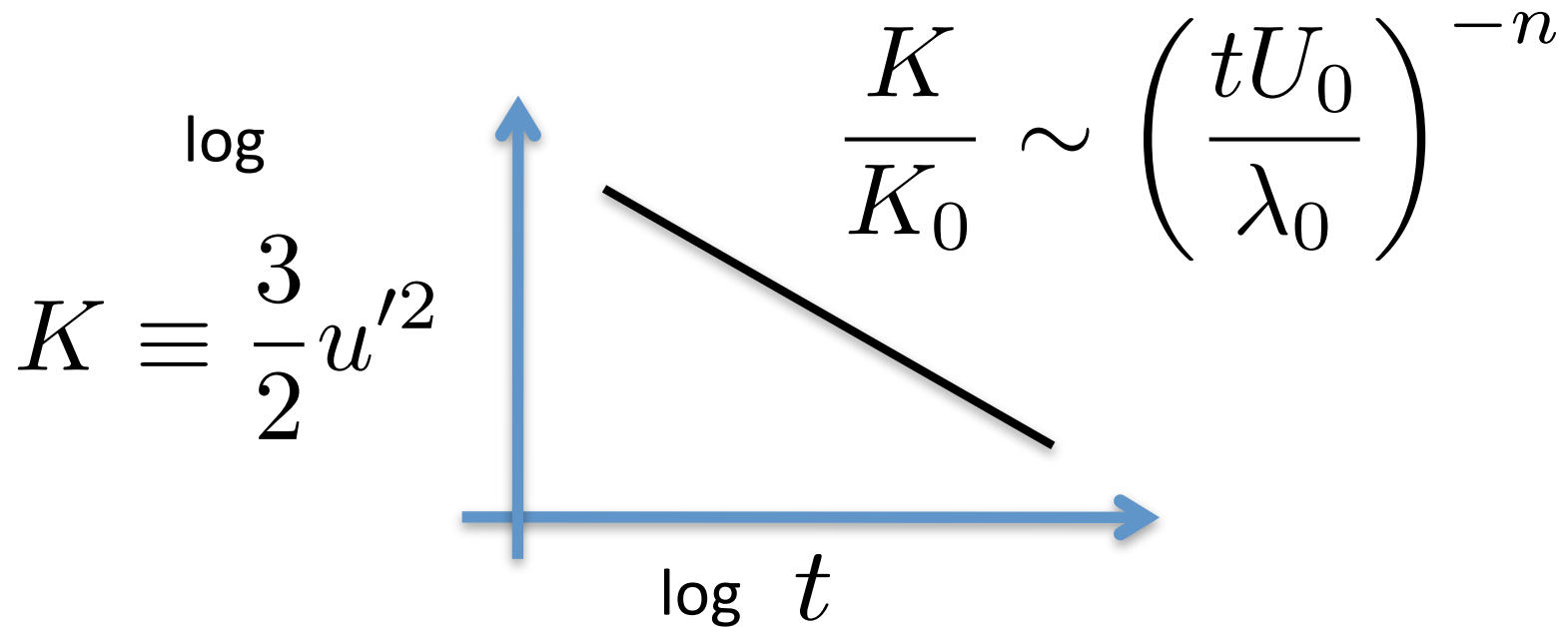
$$u_x(\mathbf{x}, t_0) = U_0 \sin\left(\frac{y}{\lambda_y}\right)$$

$$\partial_t \mathbf{u} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla P + \nu \nabla^2 \mathbf{u}$$

$$\partial_t u_x = -\frac{\nu}{\lambda_y^2} u_x$$

$$u_x = U_0 e^{-\nu t / \lambda_y^2} \sin\left(\frac{y}{\lambda_y}\right)$$

Observation:



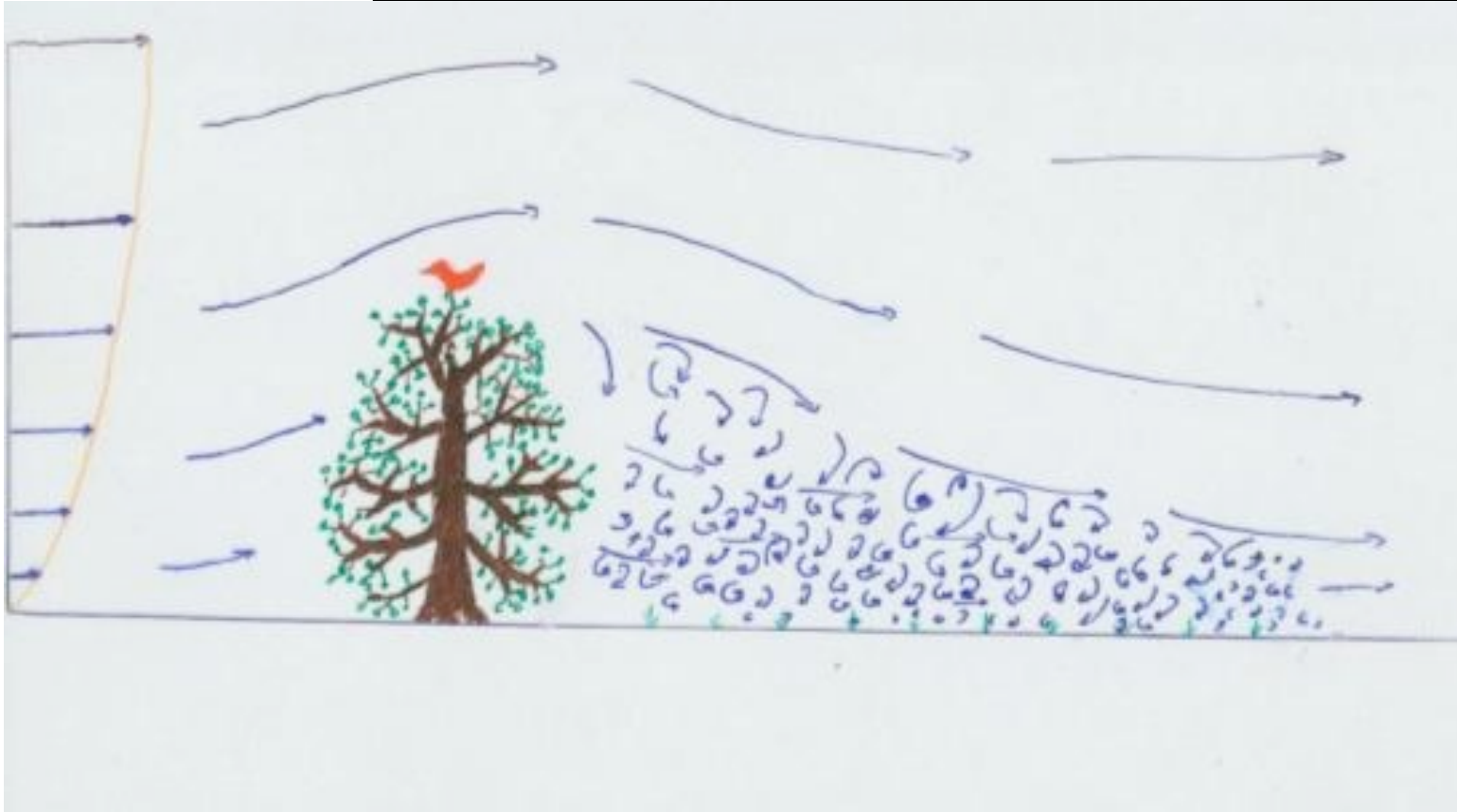
Von Kármán and Howarth, Kolmogorov, Dryden, Batchelor, Saffman, etc...



http://ict-aeolus.eu/images/horns_rev.jpg



<http://www.fao.org/docrep/010/ag127e/ag127e08.htm>



effect of Reynolds number on:

1. Decay of turbulence
2. Scaling in turbulence

RATE OF DECAY

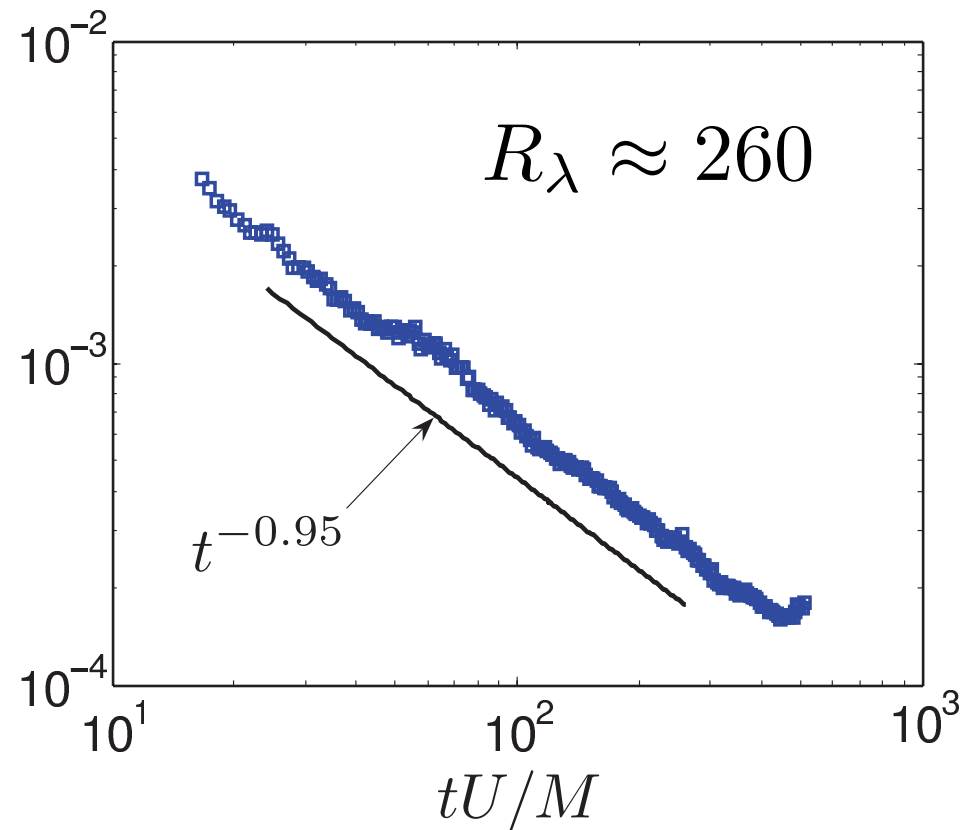
1. $r, t \Rightarrow r/L(t)$

2. $Re = const$

$$K \sim t^{-1}$$

$$2 \frac{K}{U^2}$$

KINETIC ENERGY DECAY
AFTER MIXING LIQUID HELIUM WITH A GRID



Dryden (1941) *Q. Appl. Maths*
Speziale and Bernard (1992) *J. Fluid Mech.*

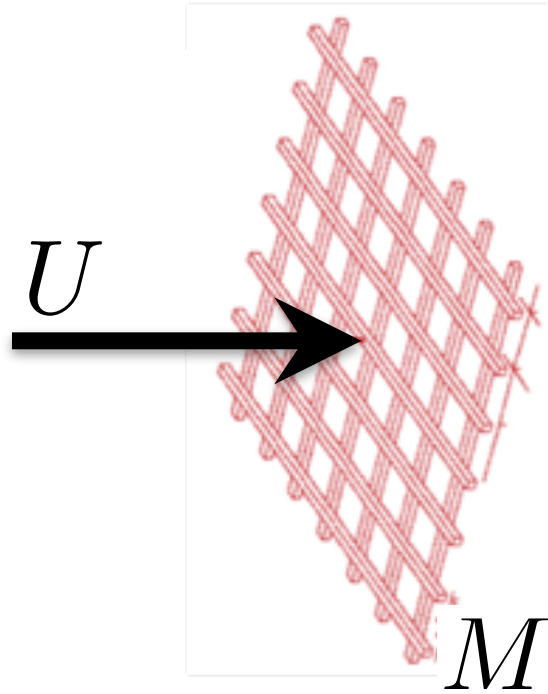
Bewley et al. (2007) *Phys. Fluids*

RATE OF DECAY

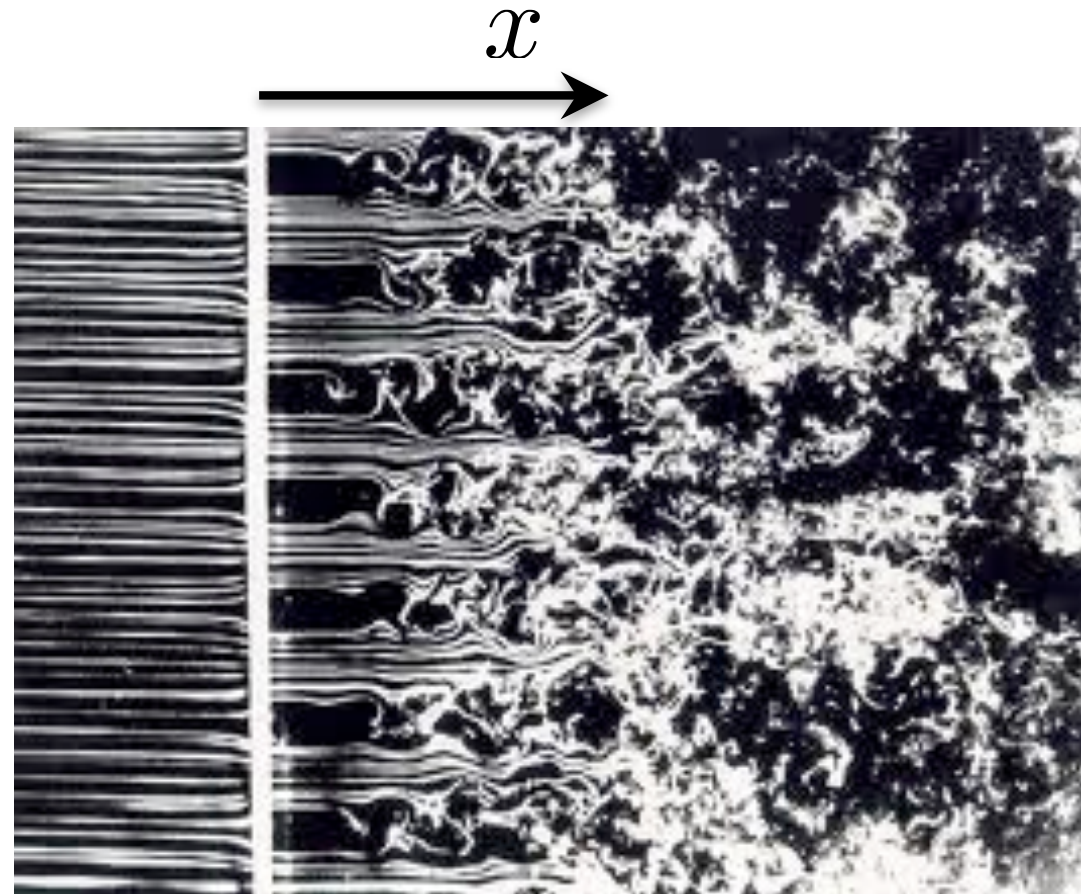
$$f(r, t) = \frac{\langle u(\vec{x}, t) u(\vec{x} + \vec{r}, t) \rangle}{u'^2}$$

$$f(r, t) \sim r^{-2} \quad \Leftrightarrow \quad K \sim t^{-6/5} \quad (\text{Saffman})$$

$$f(r, t) \sim r^{-6} \quad \Leftrightarrow \quad K \sim t^{-10/7} \quad (\text{Kolmogorov})$$



GRID TURBULENCE

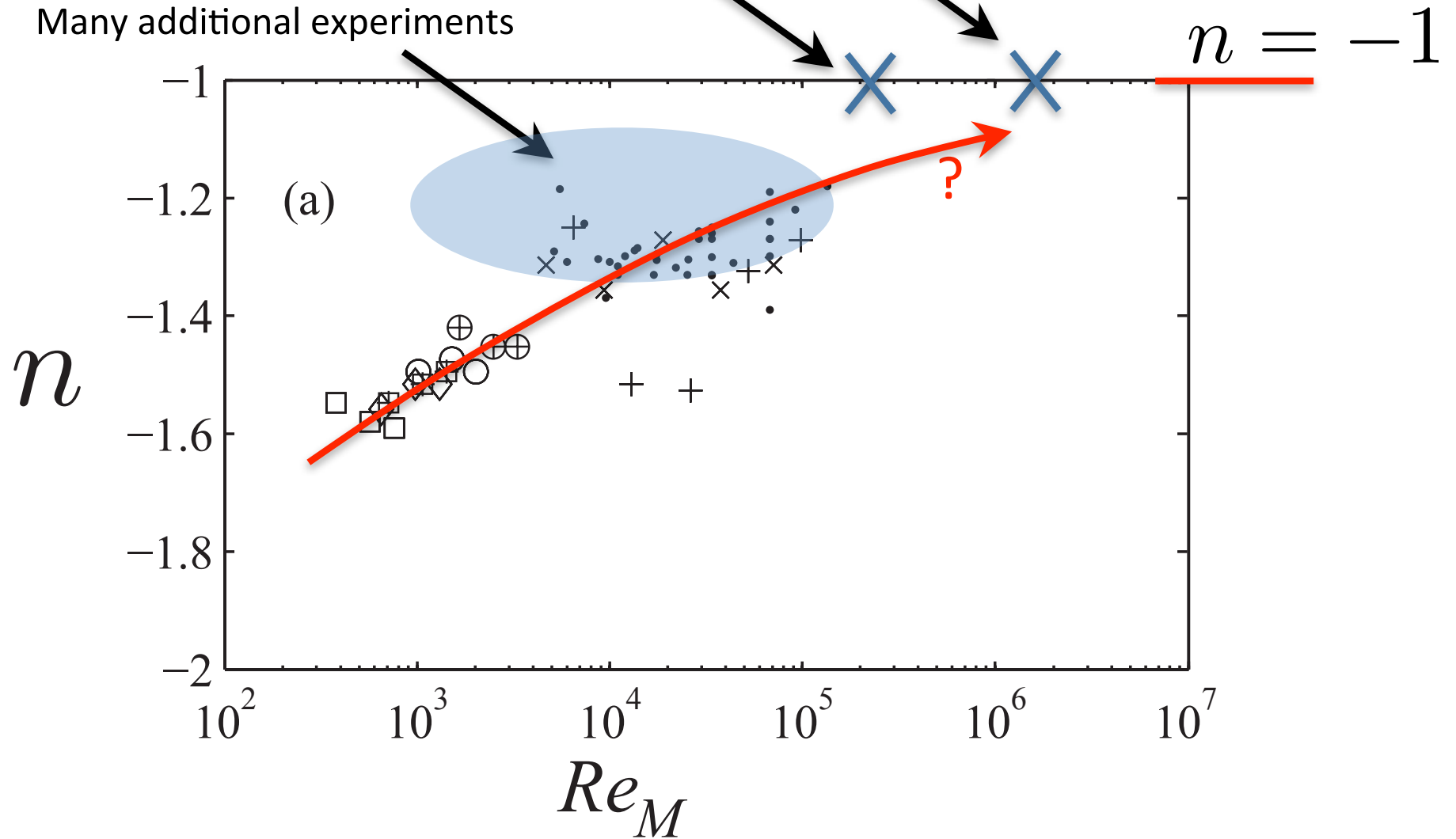


$$Re_M = \frac{UM}{\nu}$$

Kistler and Vrebalovich (1966) *J. Fluid Mech.*

Bewley et al. (2007) *Phys. Fluids*

Many additional experiments

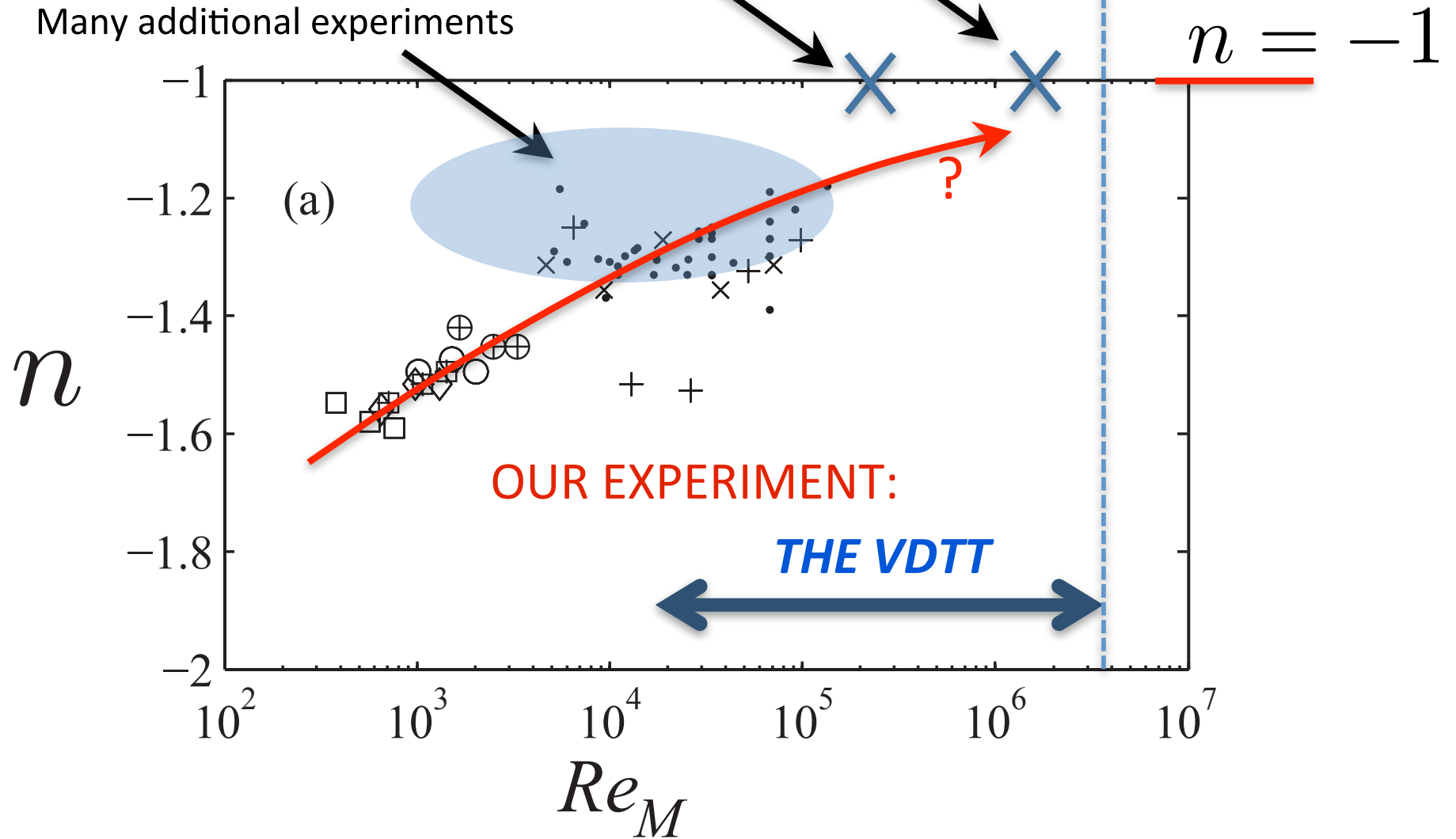


Kurian and Fransson (2009) *Fluid Dyn. Res.*

Kistler and Vrebalovich (1966) *J. Fluid Mech.*

Bewley et al. (2007) *Phys. Fluids*

Many additional experiments



Kurian and Fransson (2009) *Fluid Dyn. Res.*

THE VARIABLE DENSITY TURBULENCE TUNNEL (VDTT)

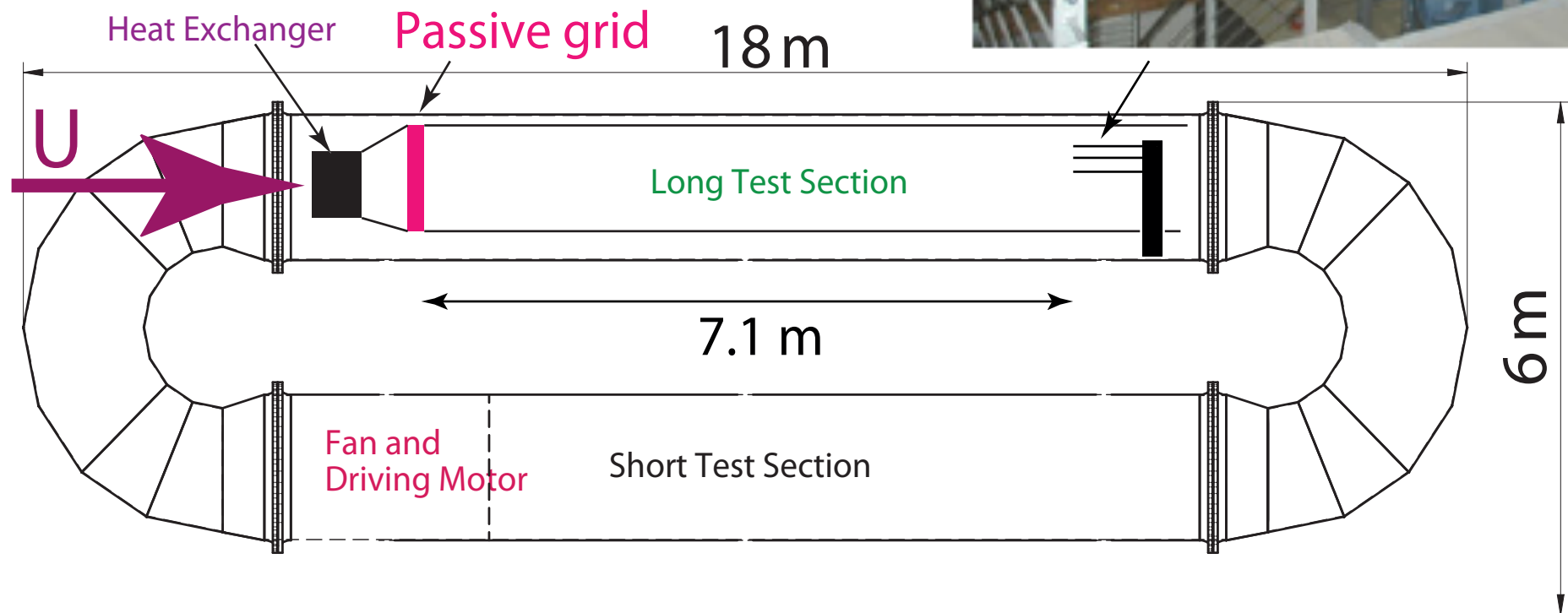


Bewley, Nobach, Sinhuber, Xu, Bodenschatz (2014) *under review*.

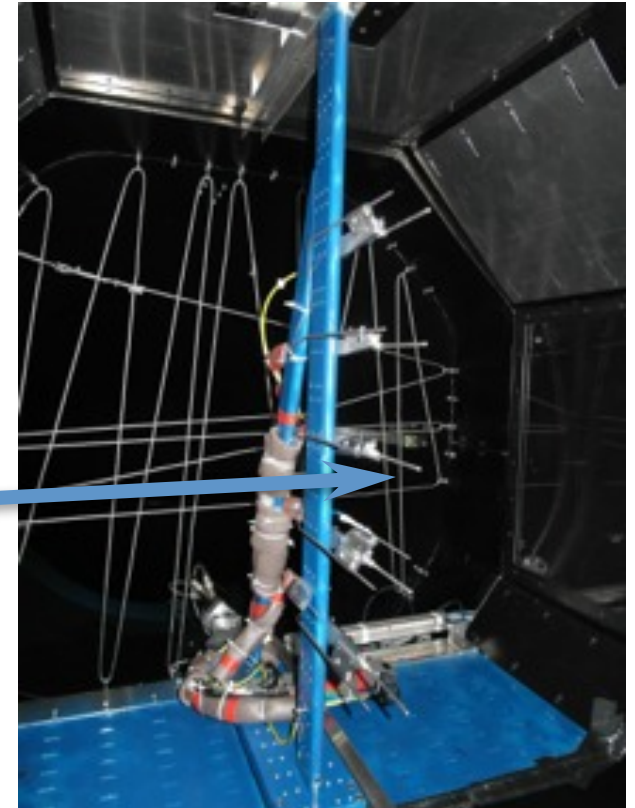
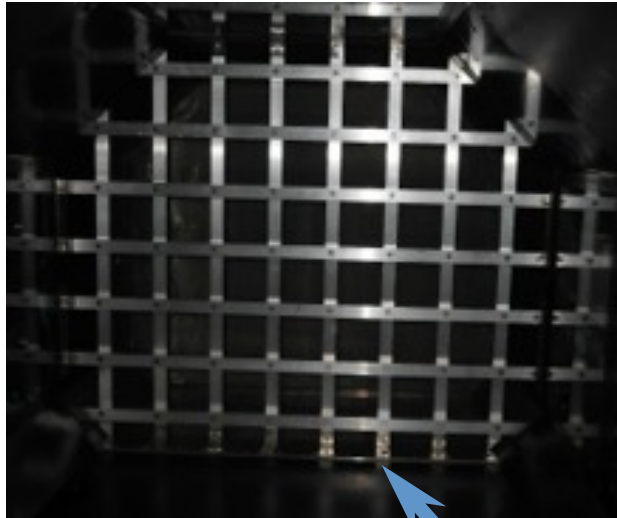
$$Re = \frac{\rho U L}{\mu}$$

Air and Sulfur Hexafluoride gas (SF_6)

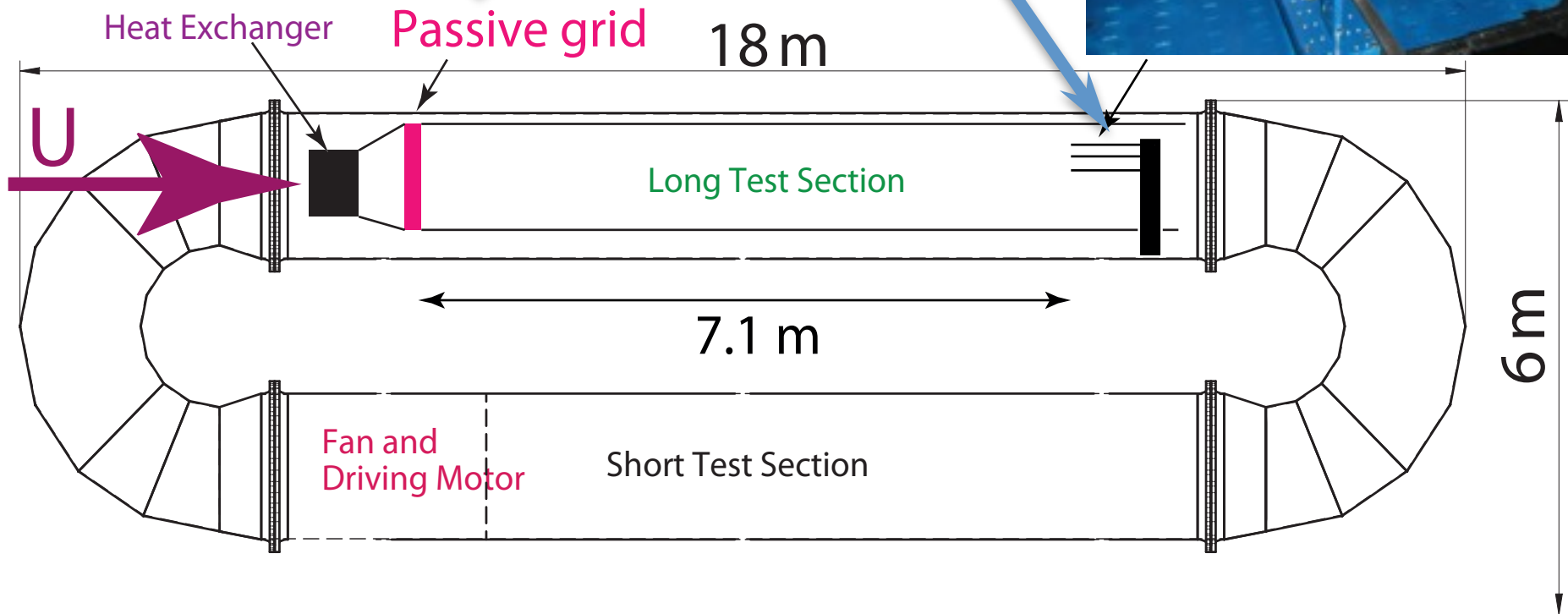
$$U \leq 5 \text{ m/s} \quad \eta \geq 20 \text{ } \mu\text{m}$$
$$\tau_\eta \geq 2 \text{ ms}$$



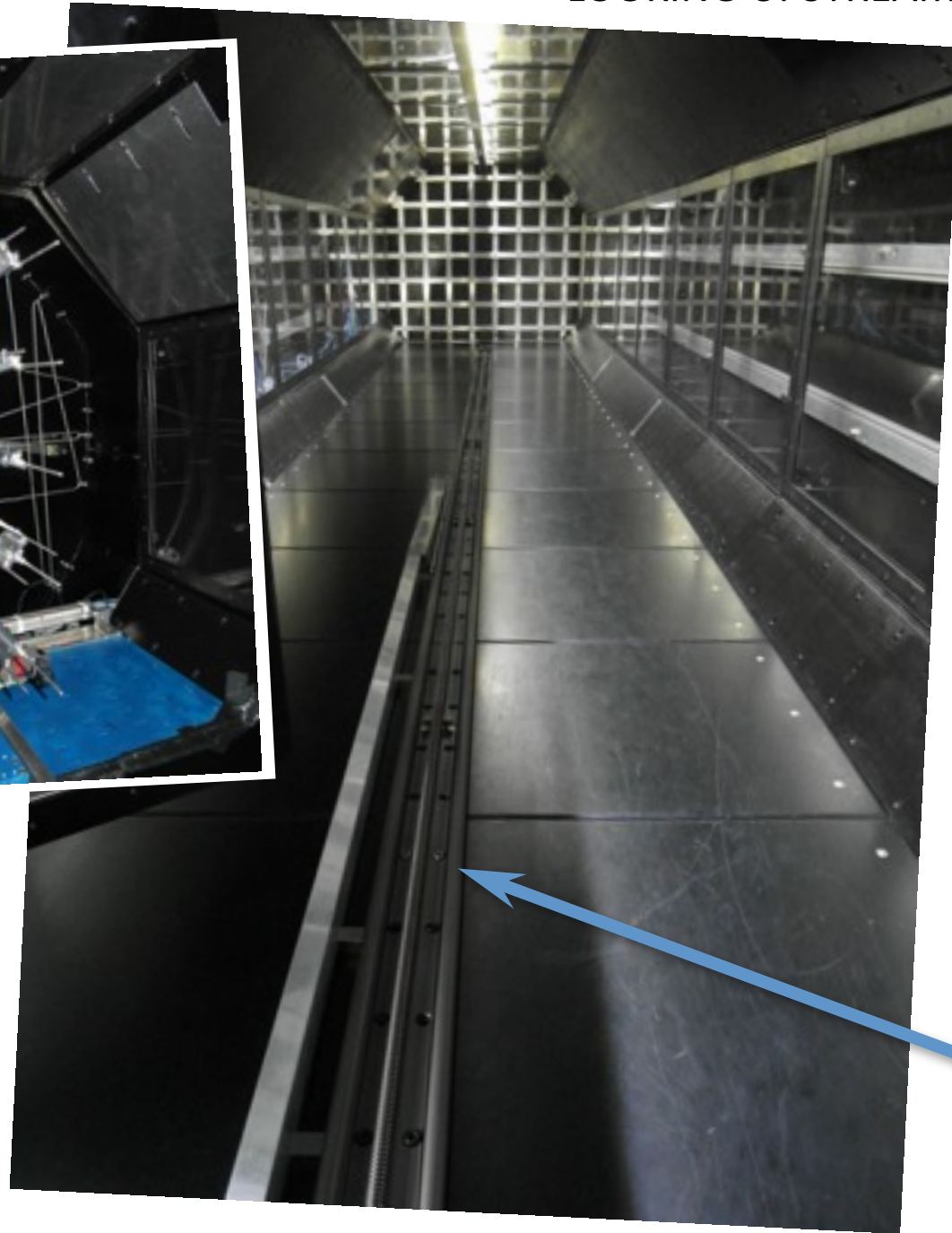
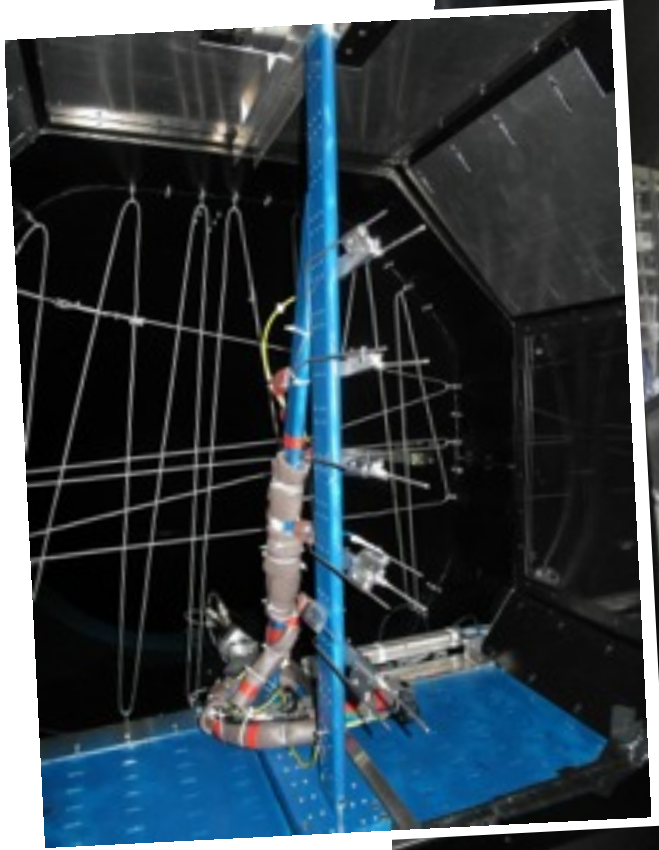
M = 180 mm



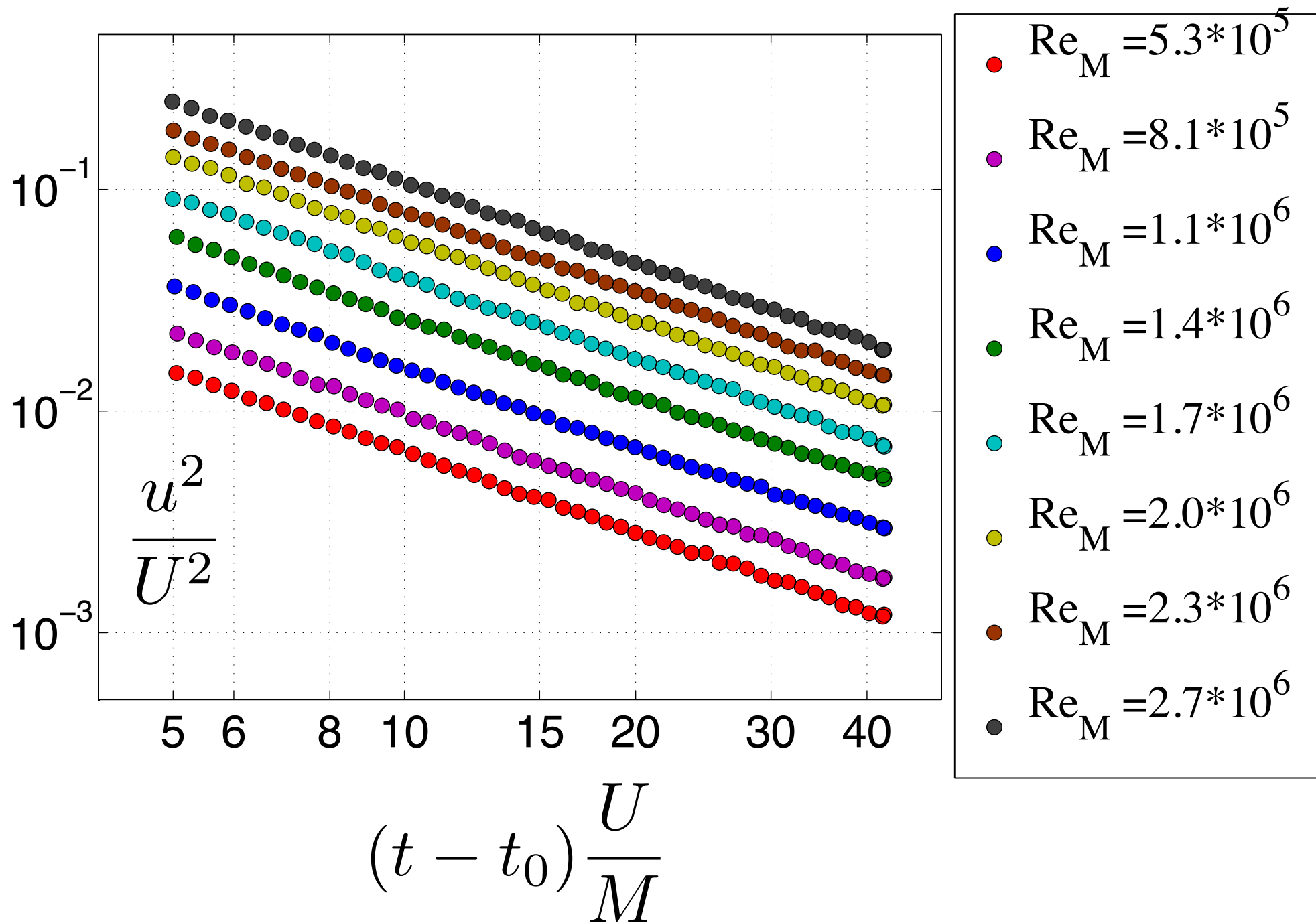
HOT WIRE
PROBES
ON TRAVERSE



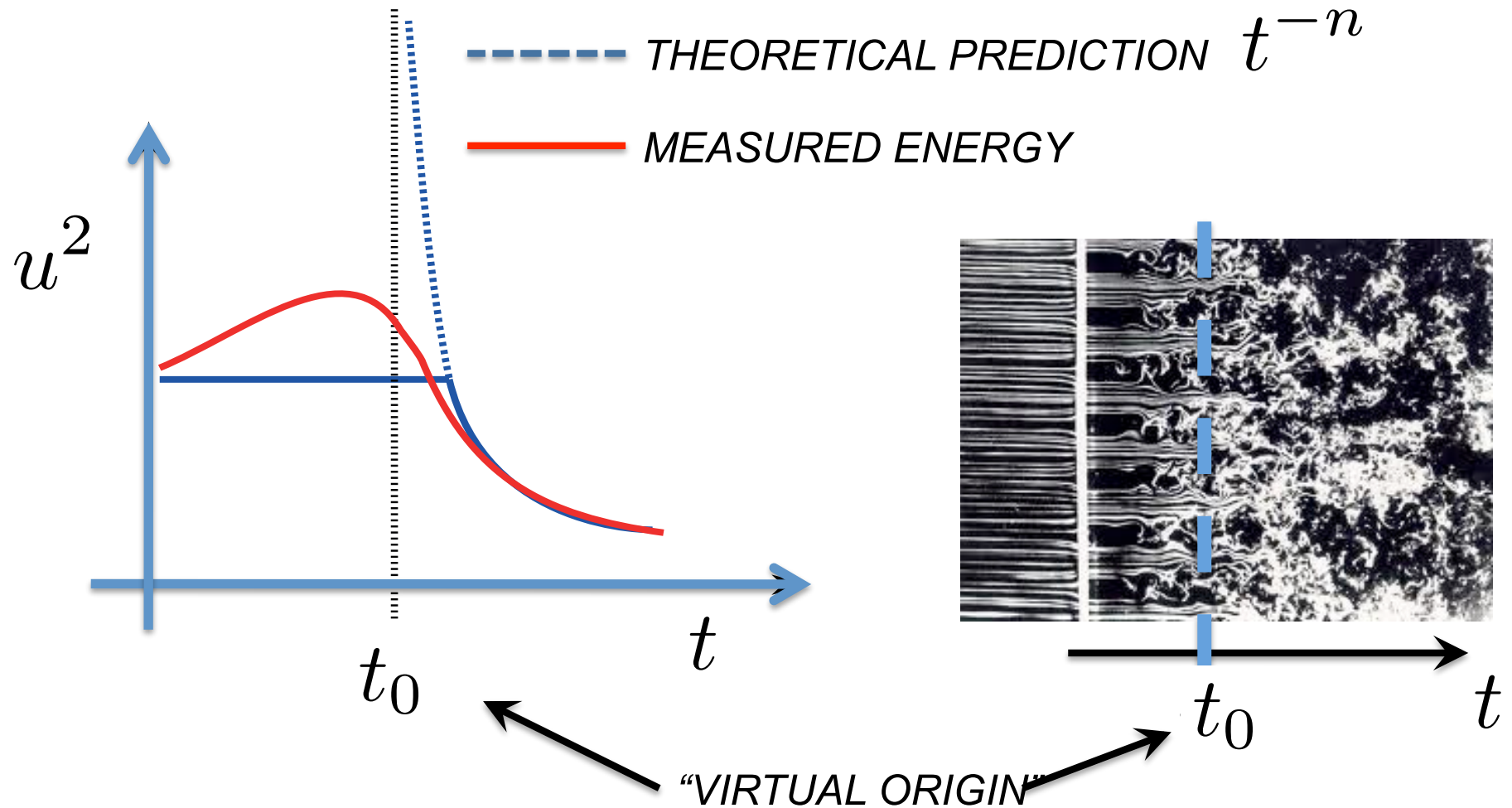
LOOKING UPSTREAM



TRAVERSE

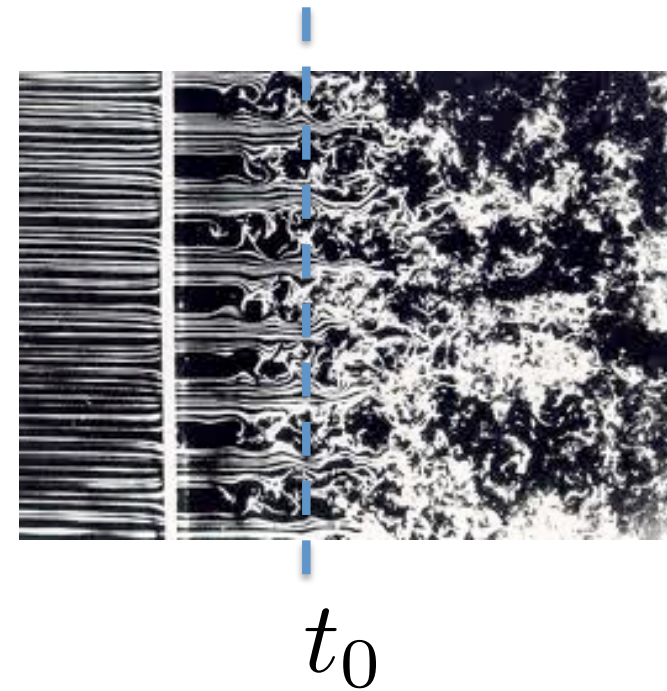
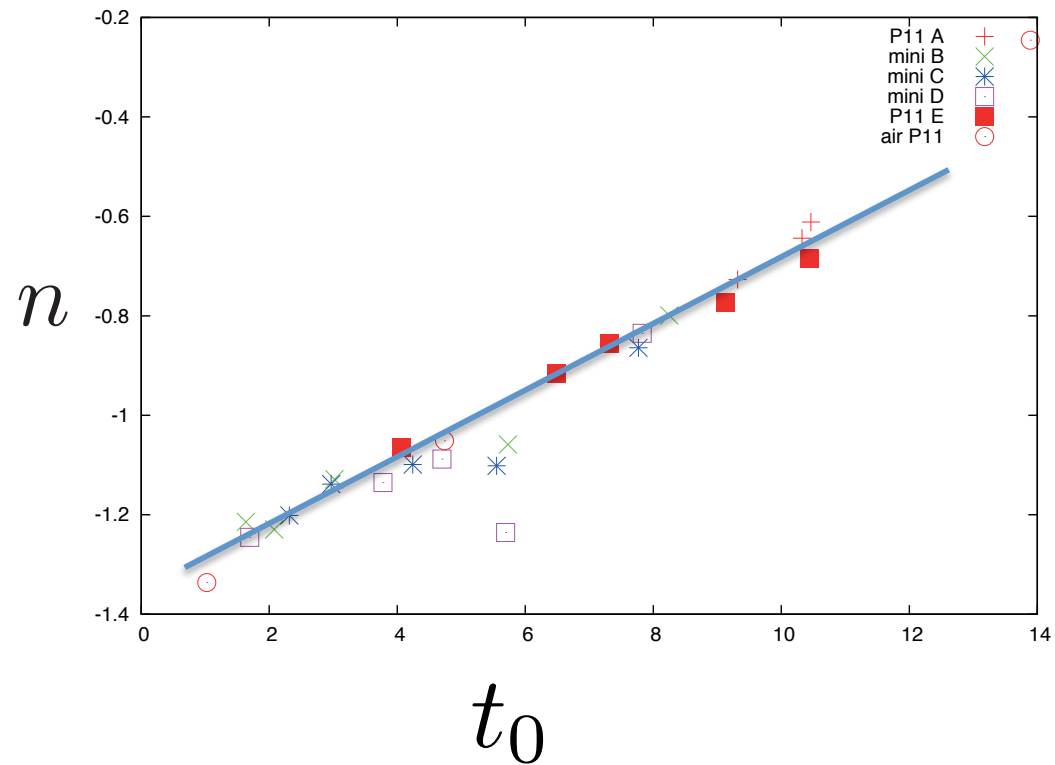


TO DETERMINE THE DECAY RATE IN AN EXPERIMENT:



least squares fit to: $u^2 = A(t - t_0)^{-n}$

$$u^2 = A(t - t_0)^{-n}$$



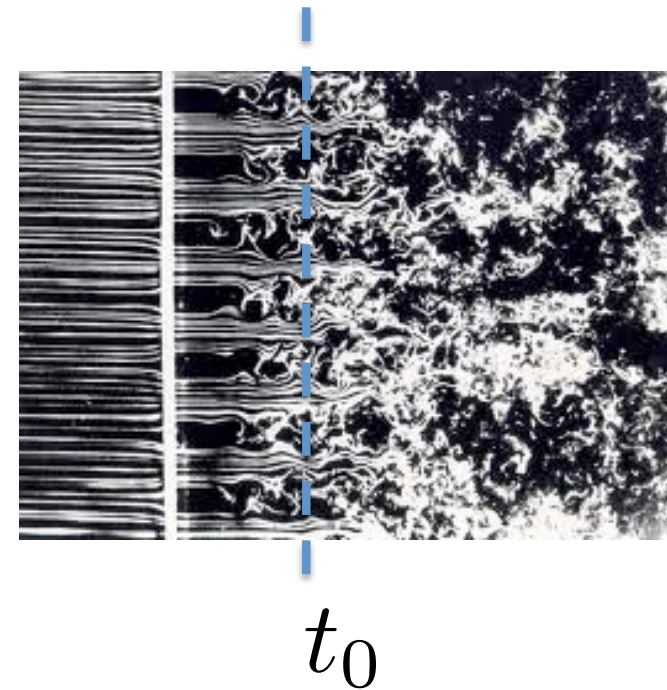
e.g. Mohammed and LaRue (1990) *J. Fluid Mech.*

$$u^2 = A(t - t_0)^{-n}$$

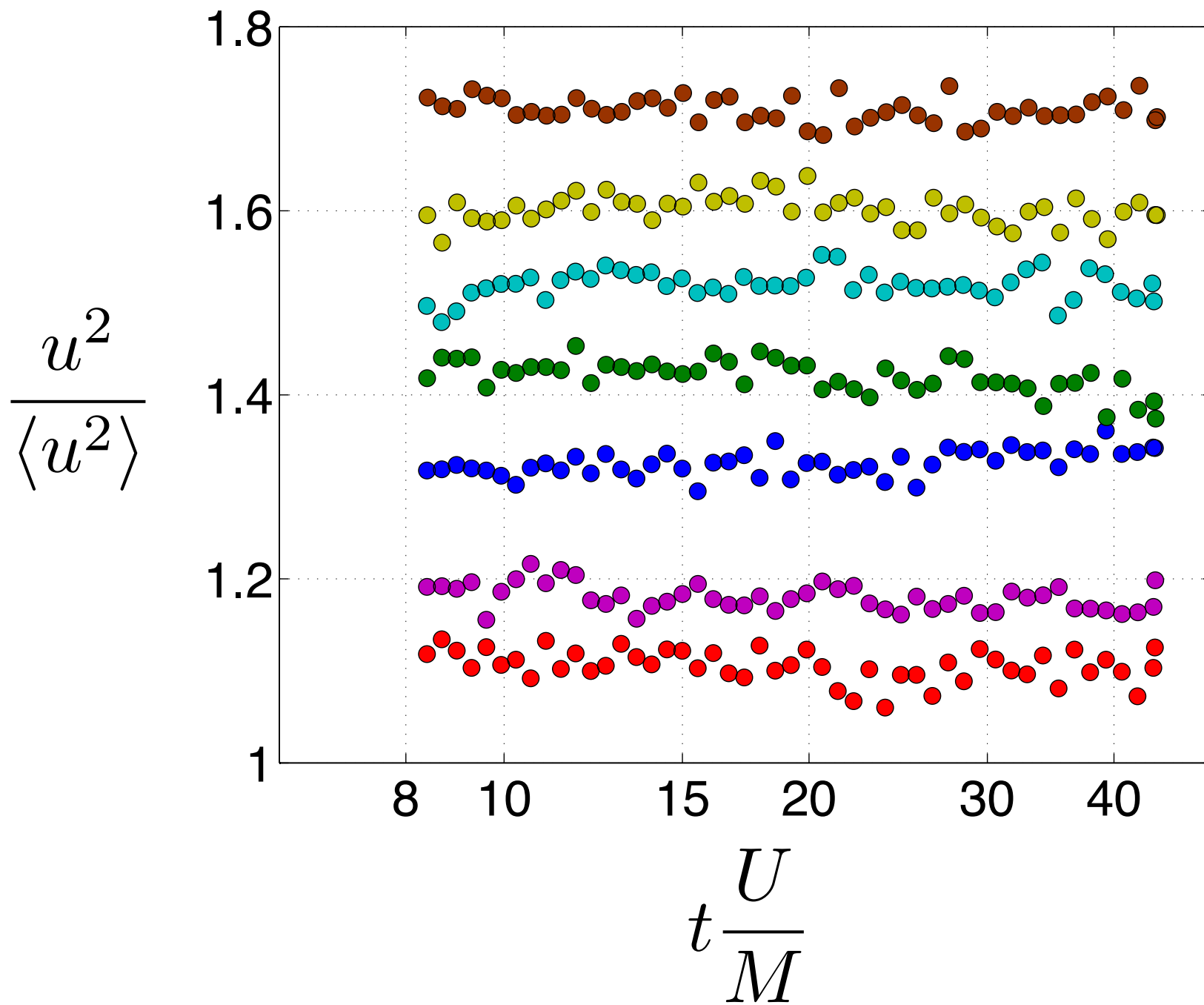
IDEA:

$$u_i^2 \sim (u_0^2)^{n_i/n_0}$$

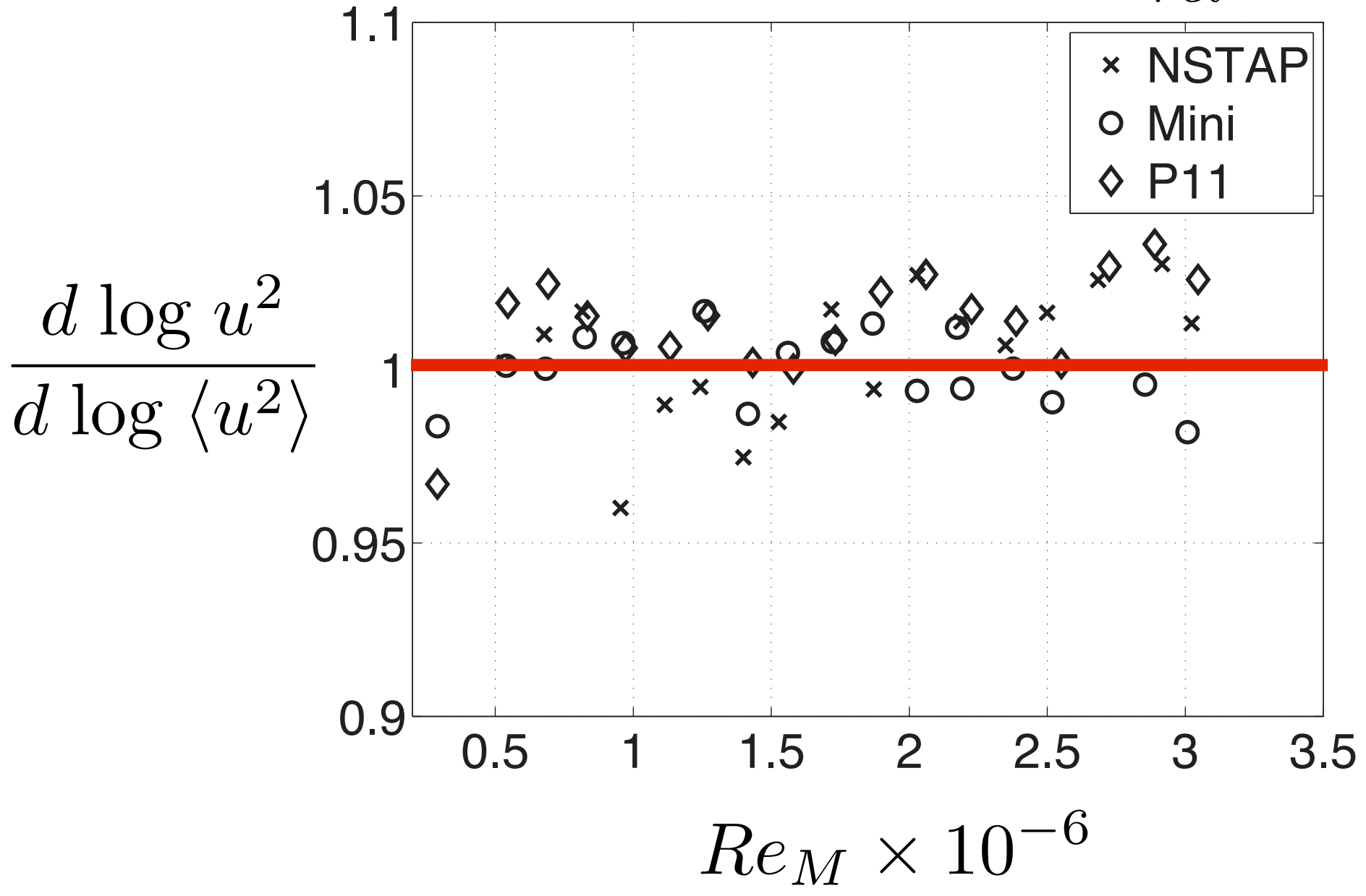
Eliminates dependence on t_0



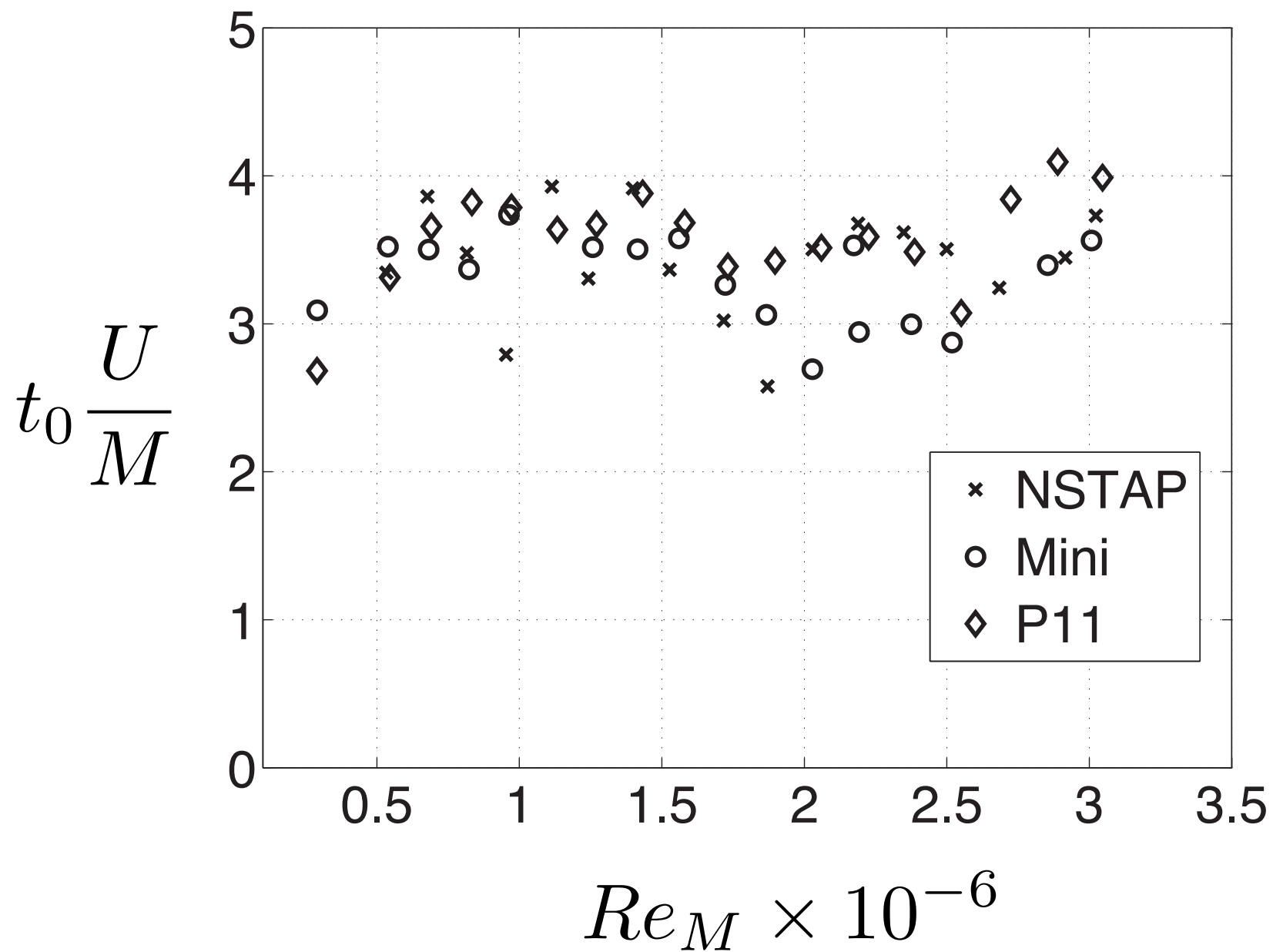
Valid if variation in virtual origin with Reynolds number is small.



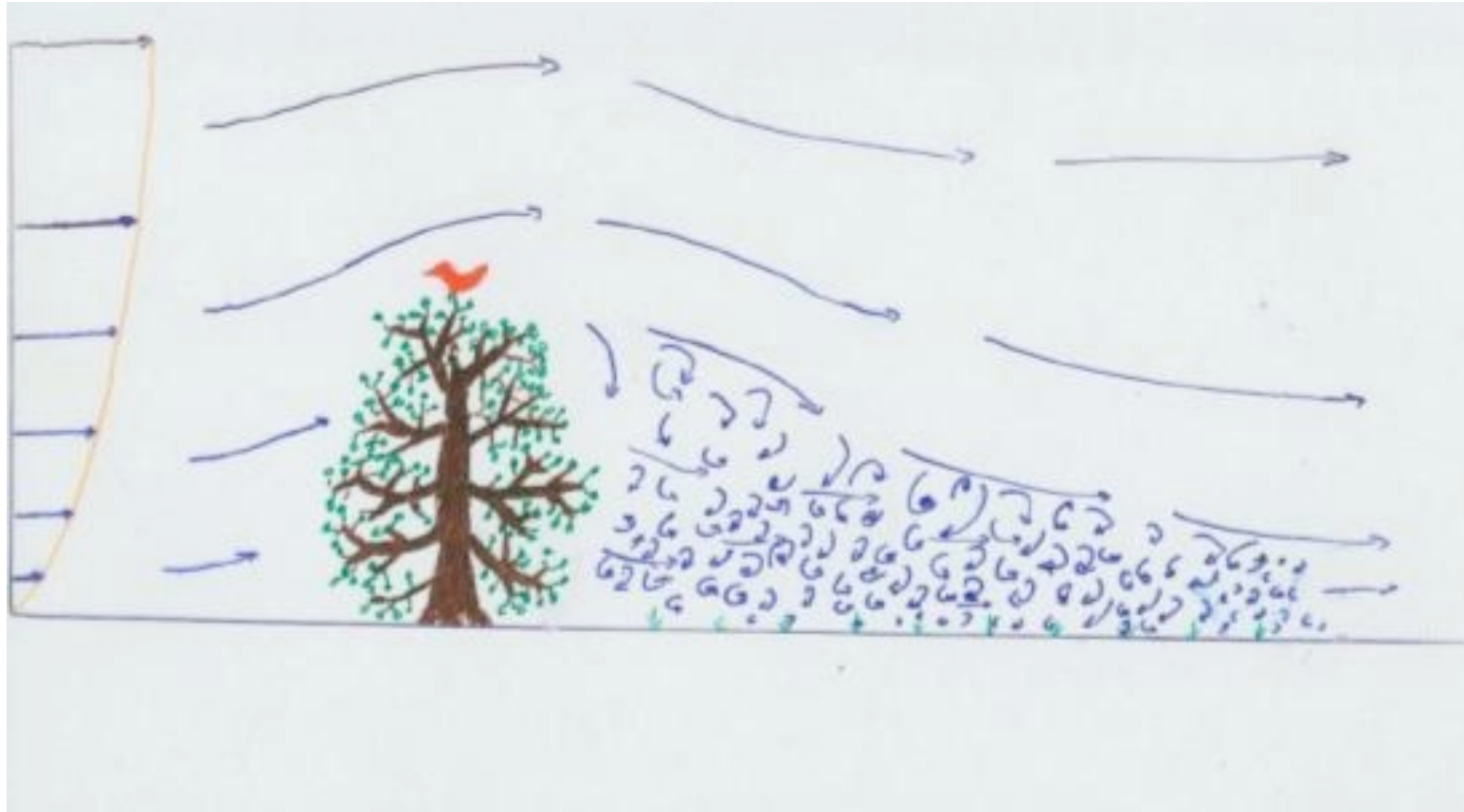
RELATIVE DECAY EXPONENT n_{rel}



VIRTUAL ORIGIN



SCALING



<http://www.fao.org/docrep/010/ag127e/ag127e08.htm>

$$\delta v = v(x + r, t) - v(x, t)$$

$$S_n(r) = \langle \delta v^n \rangle$$

SCALING

$$\frac{3}{r^3} \int_0^r \frac{\partial}{\partial t} S_2(s, t) ds + S_3 = -\frac{4}{5} \epsilon r + 6\nu \frac{\partial S_2}{\partial r}$$

$$\epsilon = \nu \left\langle \frac{\delta u_i}{\delta x_j} \frac{\delta u_i}{\delta x_j} \right\rangle$$

for *locally* isotropic turbulence
and sufficiently
high Reynolds number:

$$\frac{3}{r^3} \int_0^r \frac{\partial}{\partial t} S_2(s, t) ds + S_3 = -\frac{4}{5} \epsilon r + 6\nu \frac{\partial S_2}{\partial r}$$

Kolmogorov (1941) Dokl. Akad. Nauk. SSSR...

by extension $S_n = C_n (\epsilon r)^{n/3}$

when $\epsilon(\vec{x}, t) = \langle \epsilon \rangle$



THE VDTT ~18 meters

30 – 60 micron

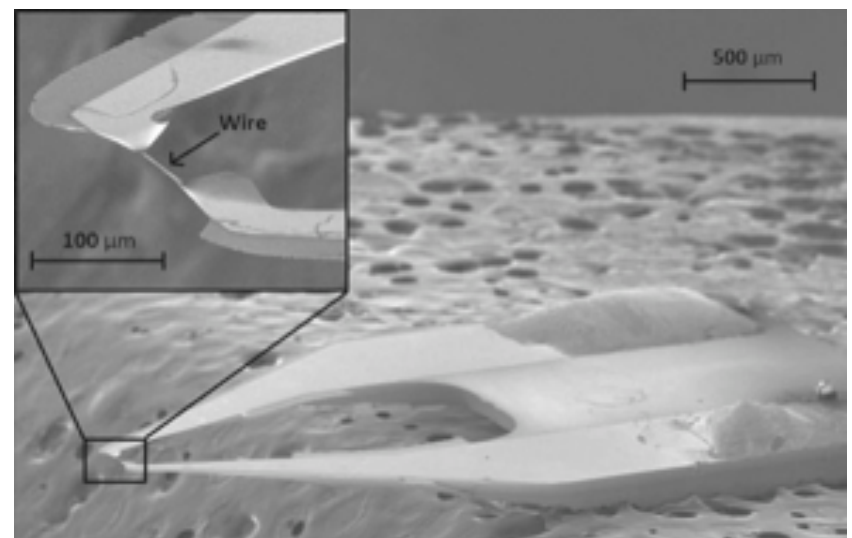
HOT WIRE PROBES

Vallikivi et al. (2011) *Expt. Fluids*

Margit Vallikivi
Marcus Hultmark
Lex Smits

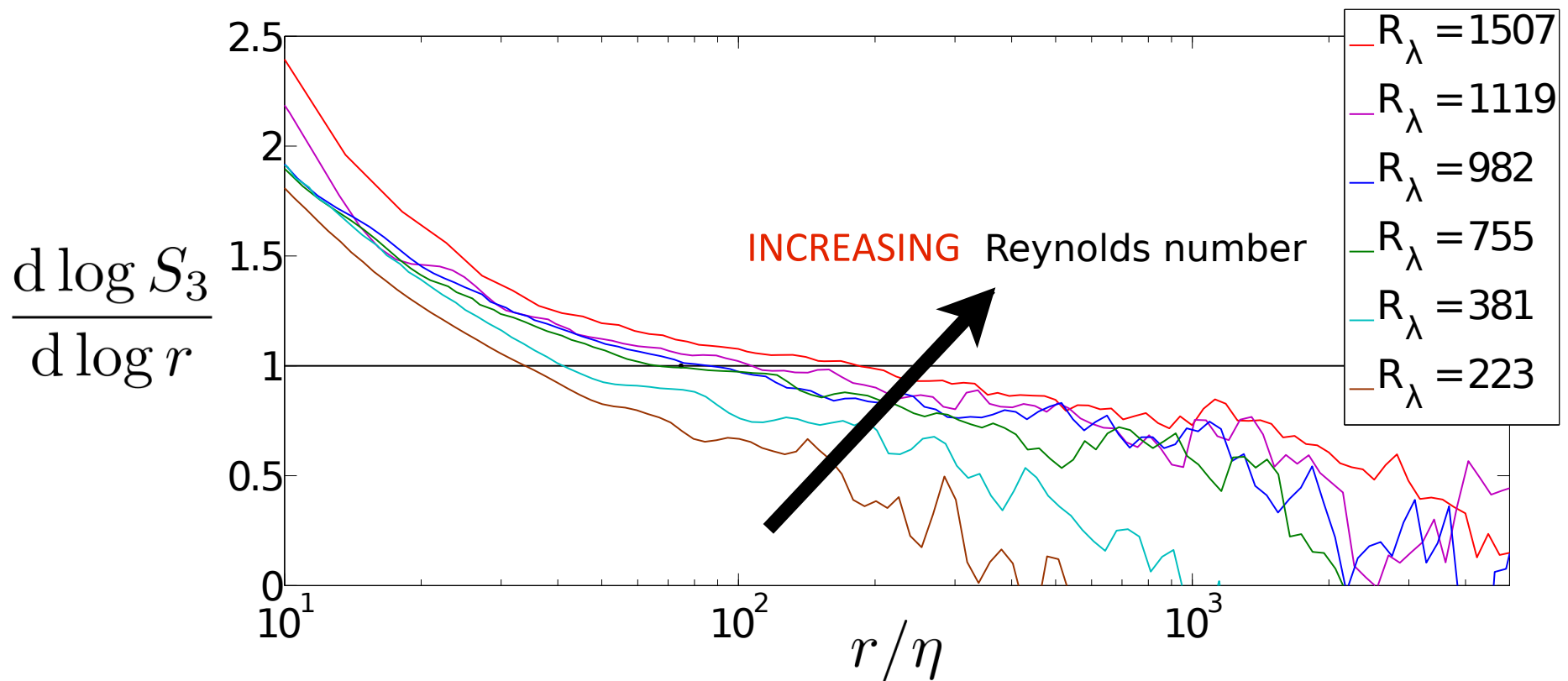
Princeton University

THE NSTAP

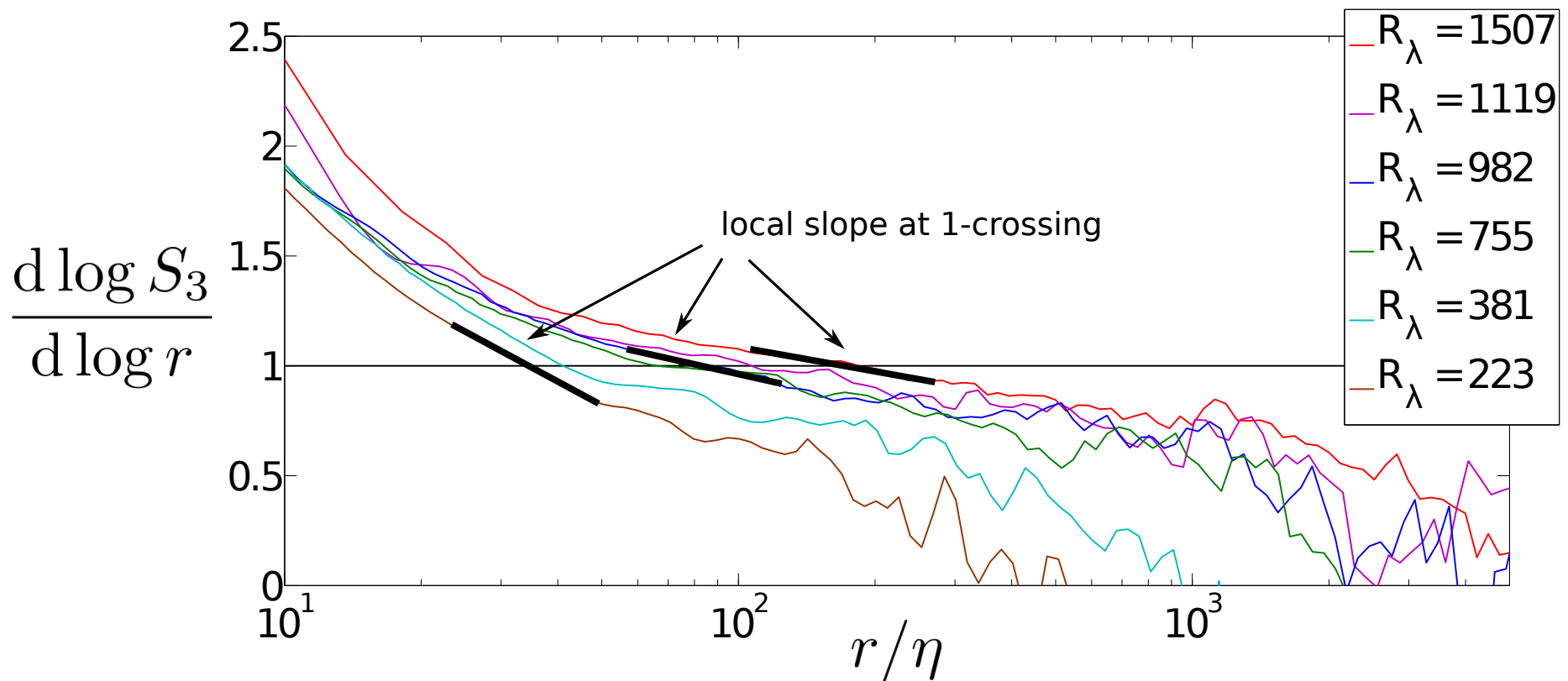


let $S_3 = Cr^{\zeta_3} \Rightarrow \frac{d \log S_3}{d \log r} = \zeta_3$

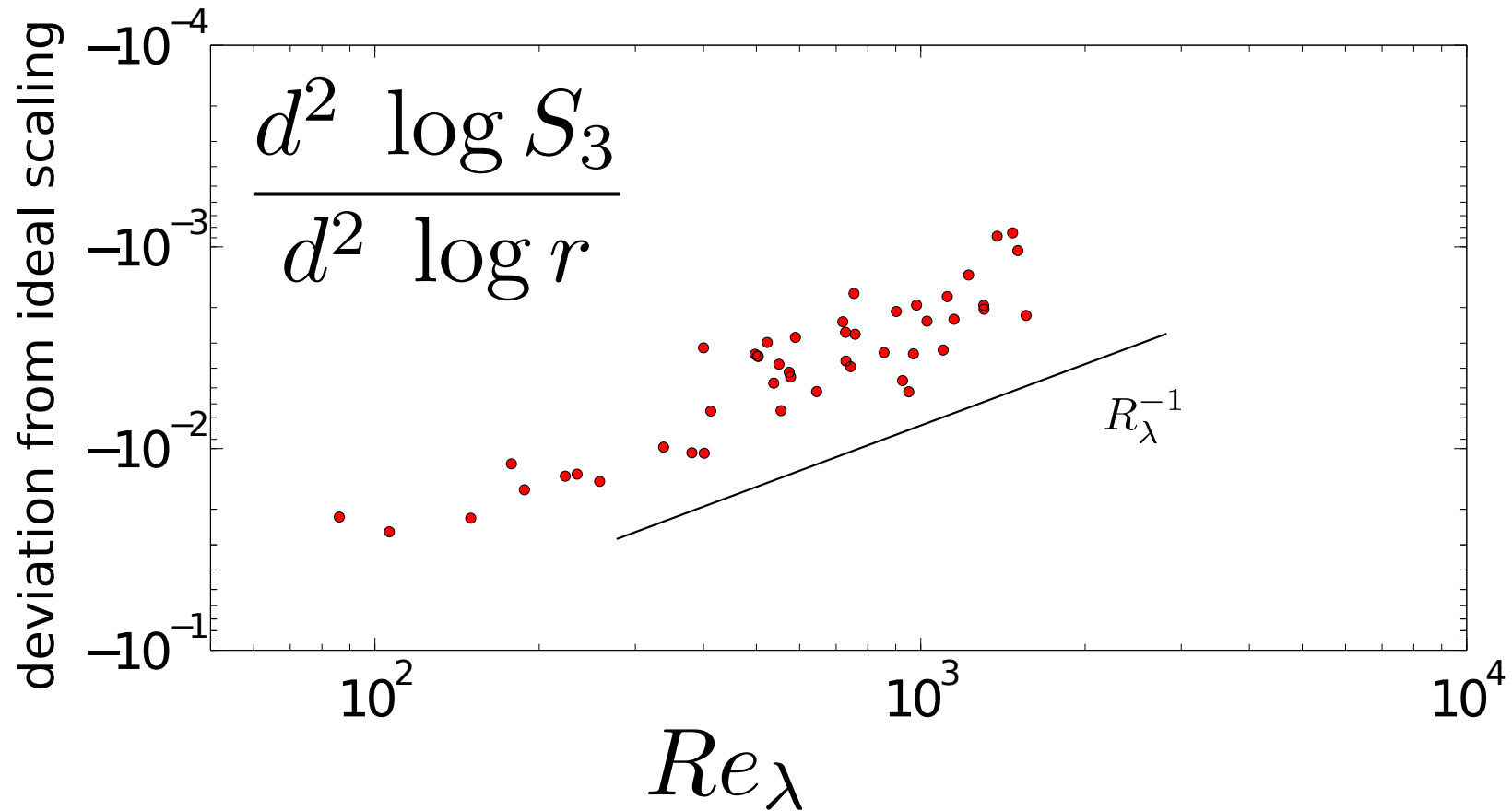
scaling would appear as a range of constant logarithmic slope



let $S_3 = Cr^{\zeta_3} \Rightarrow \frac{d \log S_3}{d \log r} = \zeta_3$



When do we first get a range of constant slope?



Qian (1999) *Phys. Rev. E*

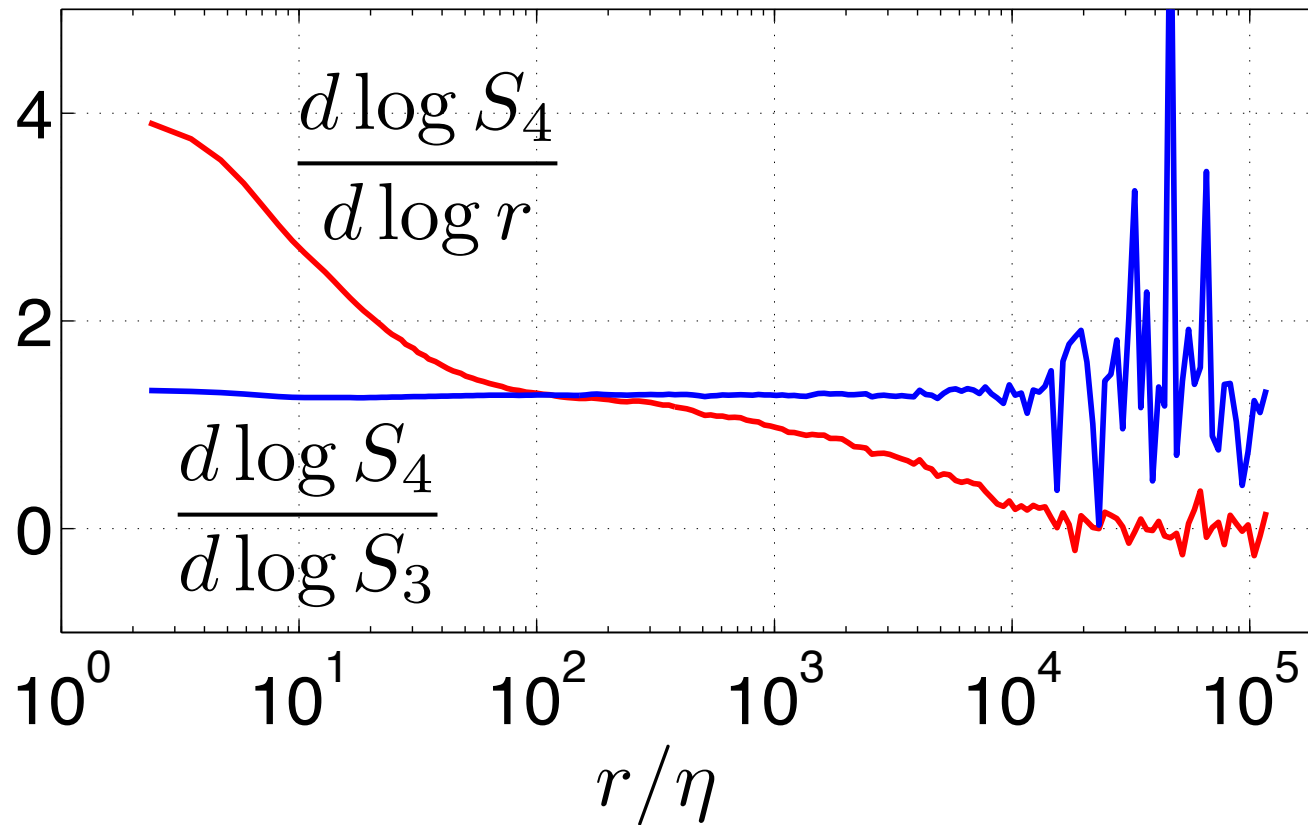
Lundgren (2002) *Phys. Fluids*

EXTENDED SELF-SIMILARITY

Benzi et al. (1993) *PRE*

$$S_n \sim S_3^{\zeta_n} \quad \zeta_n = \frac{d \log S_n}{d \log S_3}$$

$Re_\lambda = 841$

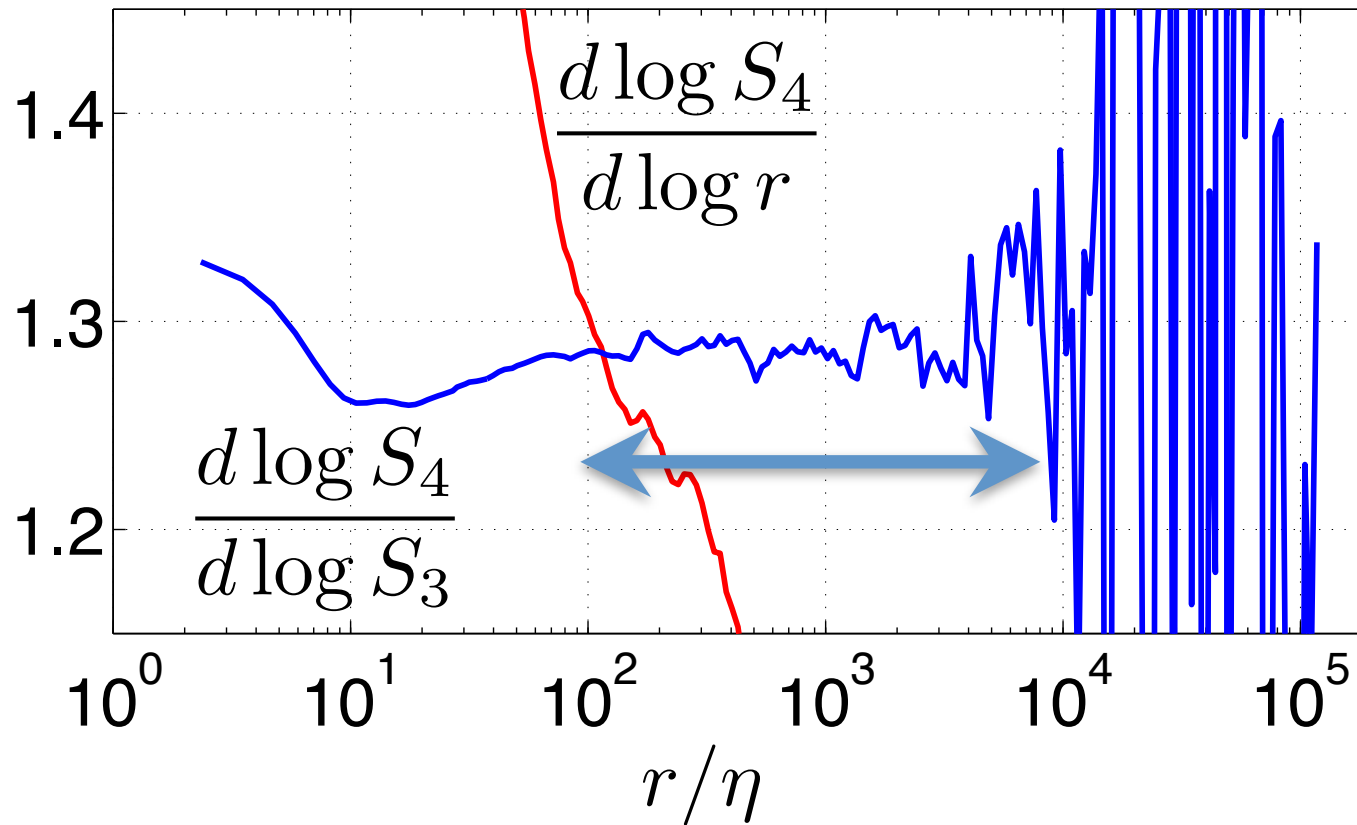


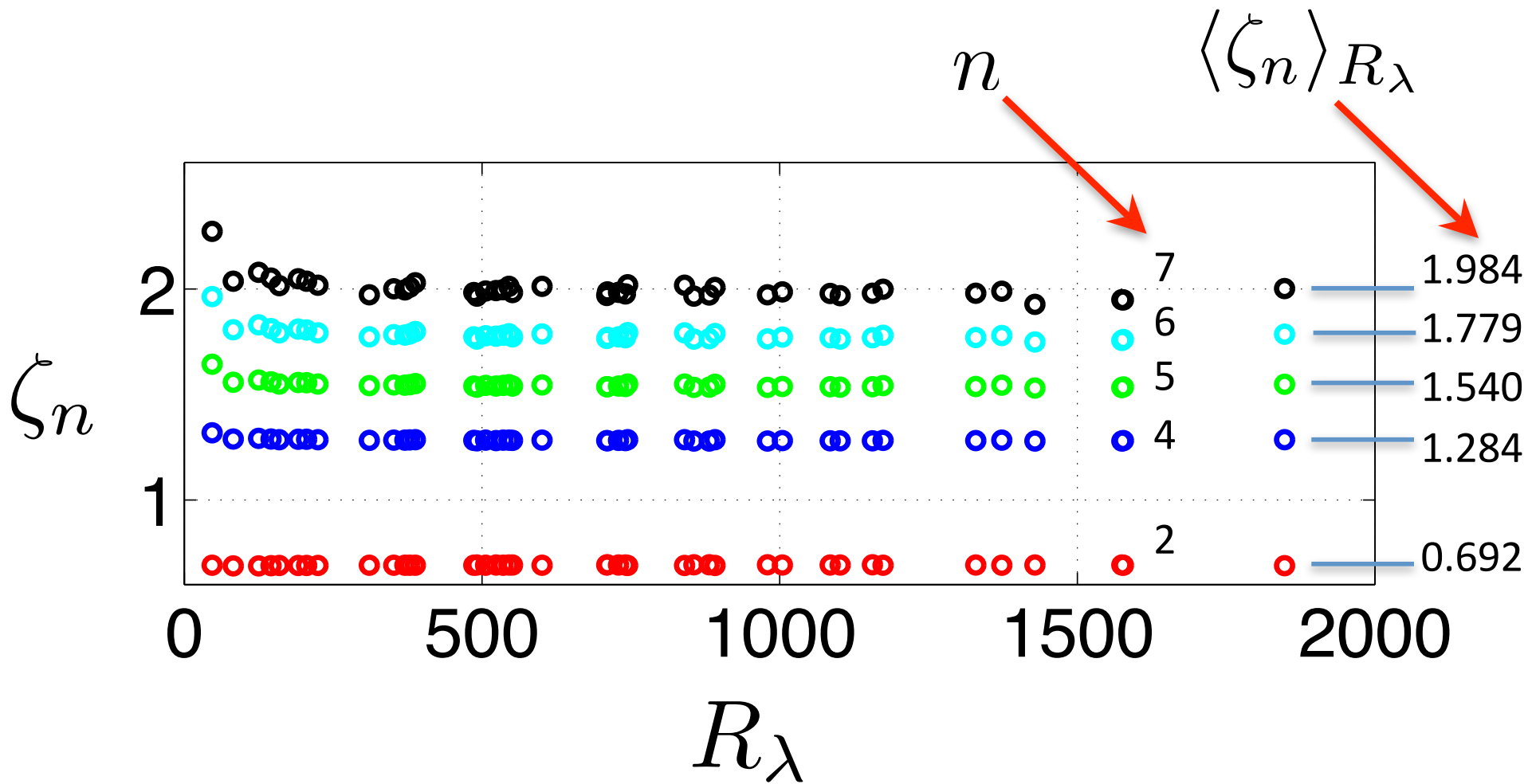
EXTENDED SELF-SIMILARITY

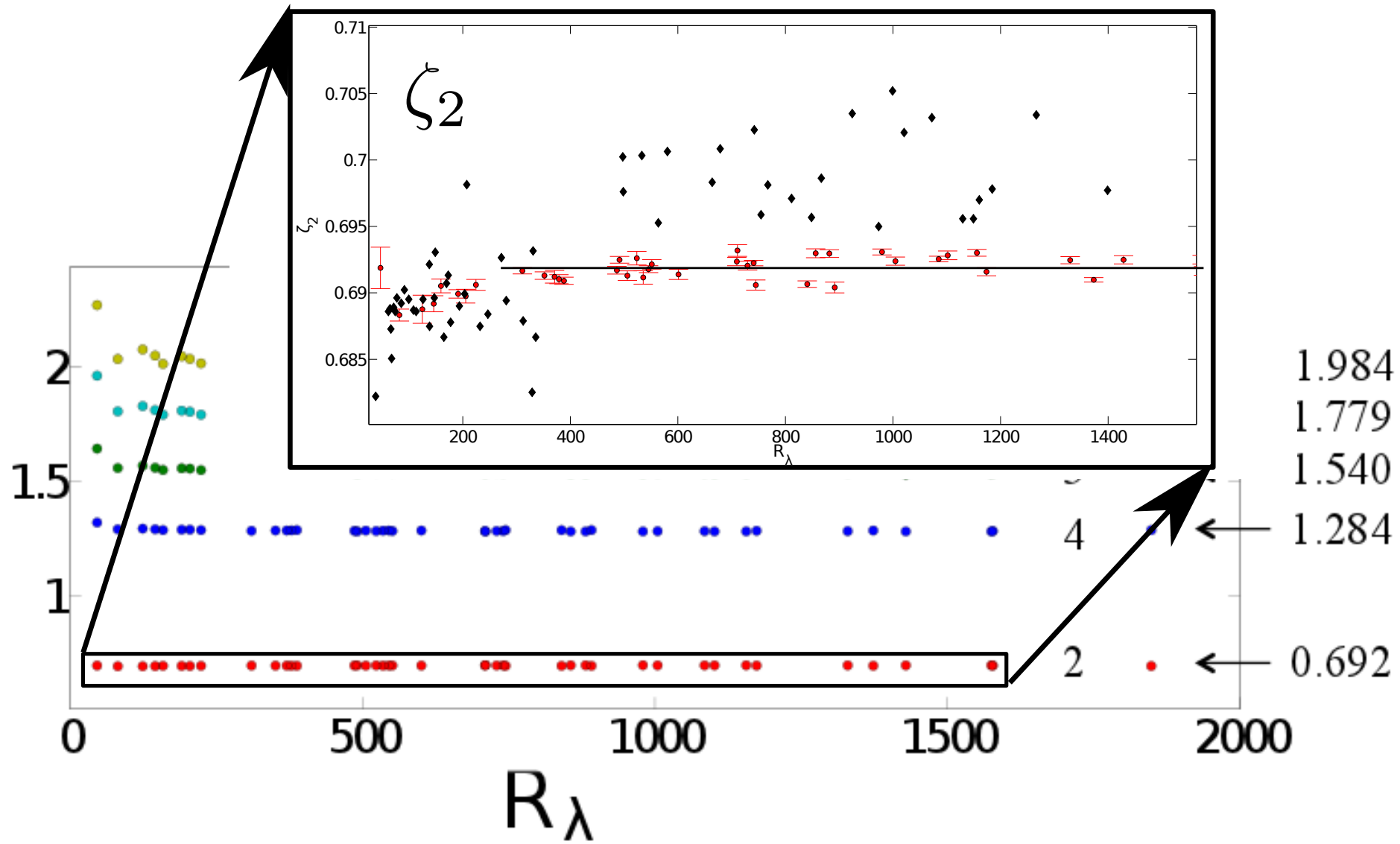
Benzi et al. (1993) *PRE*

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$Re_\lambda = 841$







Pearson and Antonia (2001) *JFM*

	VDTT	BL	DNS	S-L	K41
ζ_2	0.6915 ± 0.0006	0.708	0.699	0.695	0.666...
ζ_4	1.284 ± 0.001	1.26	1.279	1.280	1.333...
ζ_6	1.779 ± 0.009	1.71	1.772	1.778	2.000...

BL: Sreenivasan and Dhruva (1996) *Prog. Theo. Supp.*

DNS: Cao, Chen and She (1996) *PRL*

S-L: She and Lévéque (1994) *PRL*

K41: Kolmogorov (1941) *Dokl. Akad. Nauk. SSSR...*

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$$f(r, t) \sim r^{-6} \quad \Leftrightarrow \quad K \sim t^{-10/7} \quad (\text{Kolmogorov})$$

e.g. Davidson (2011) *Phys. Fluids*

Is it possible to imprint desired long-range correlations?

CONTROL OF LARGE-SCALE STRUCTURE

-for high Reynolds numbers

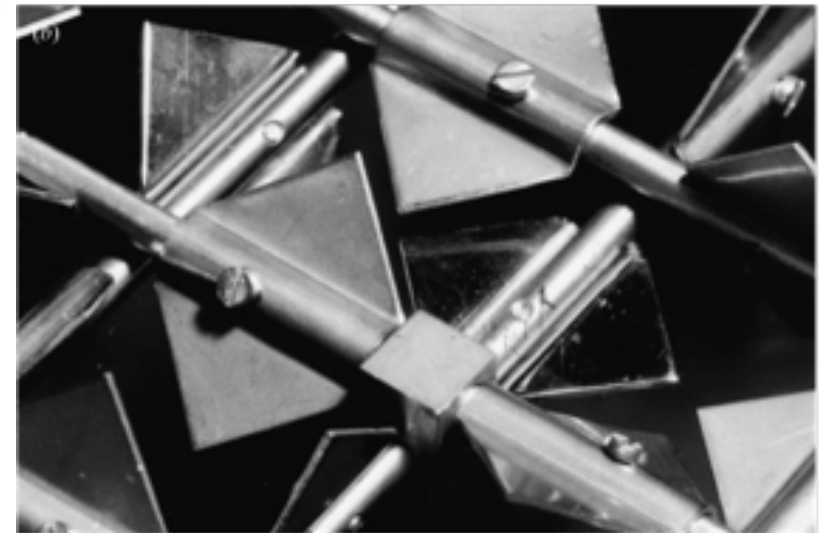
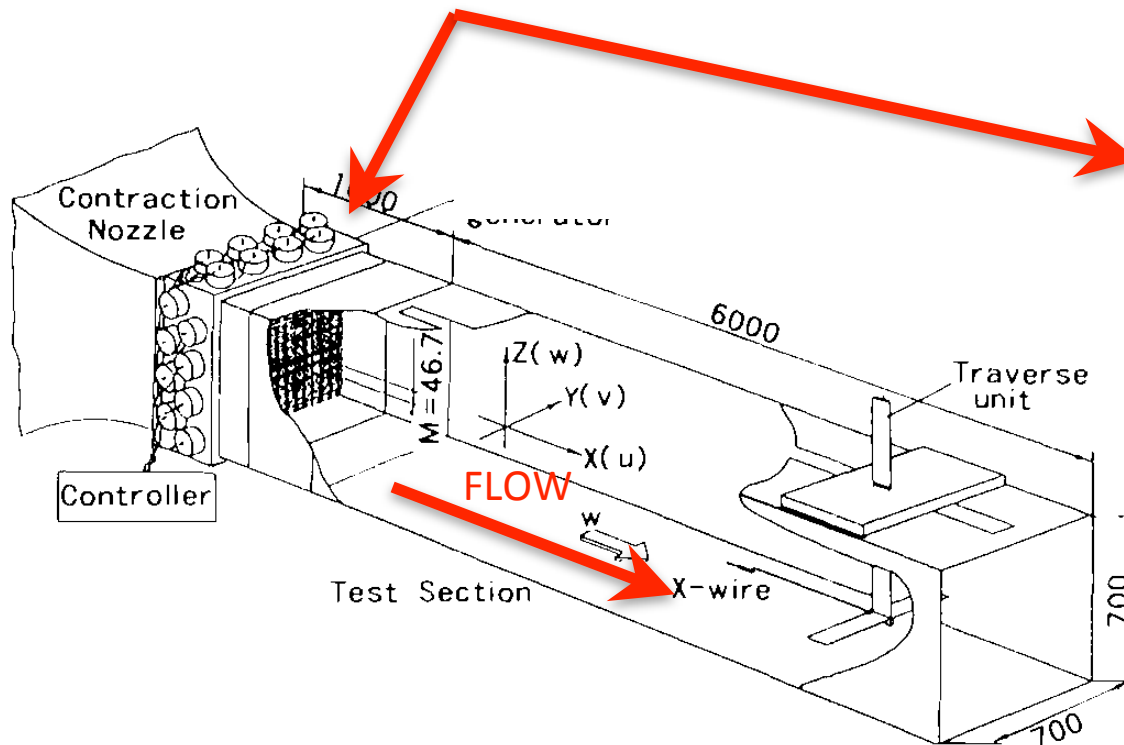
e.g. Makita (1991) *Fluid. Dyn. Res.*

-for control

e.g. Poorte and Biesheuvel (2002) *JFM*

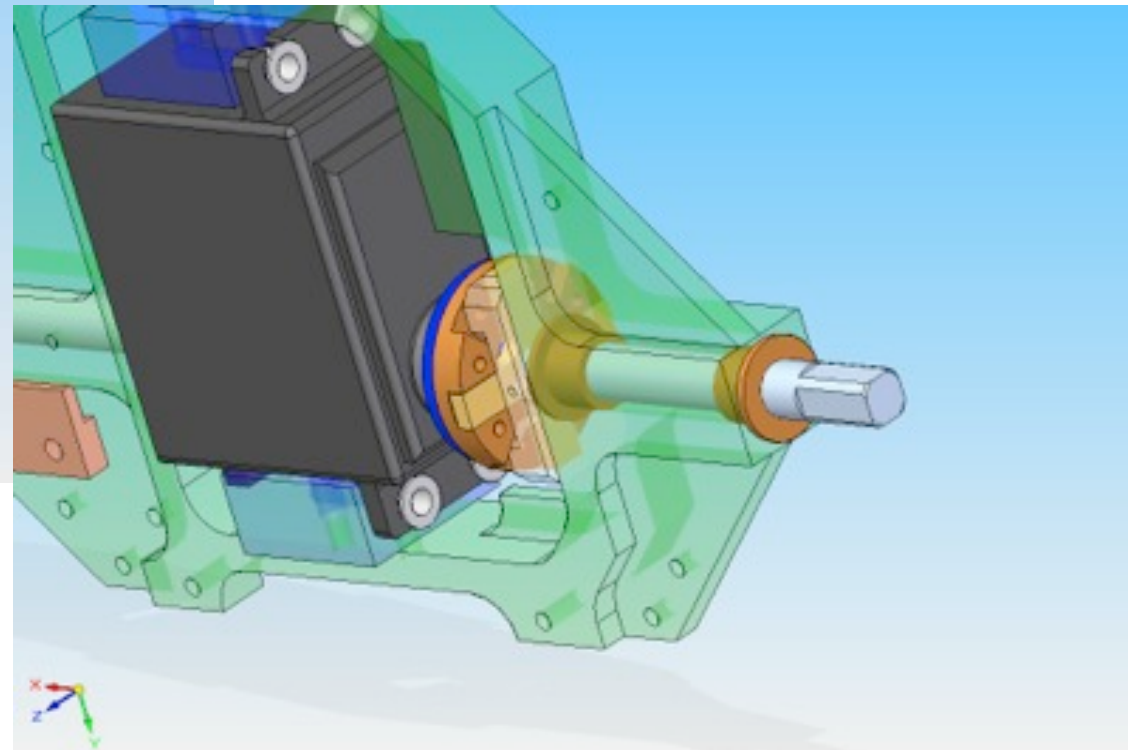
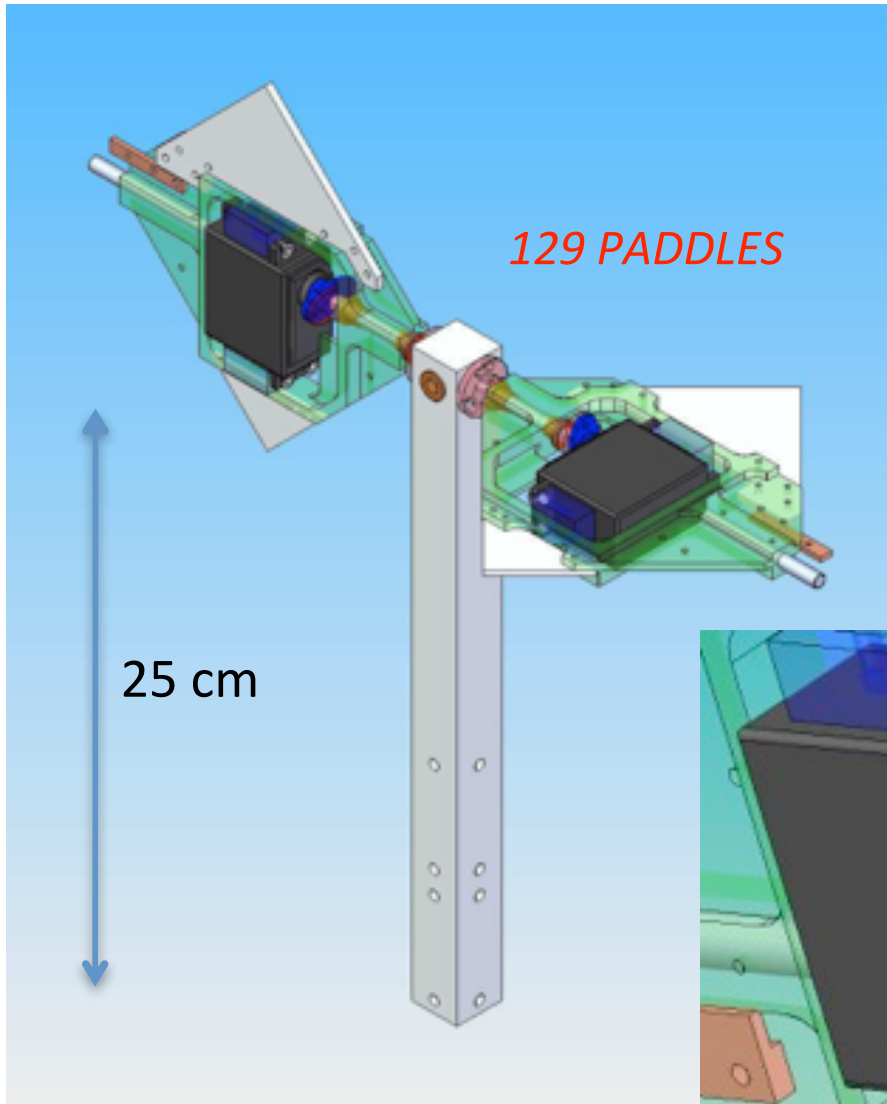
Cekli, Tipton and van de Water (2010) *PRL*

AN **ACTIVE GRID**

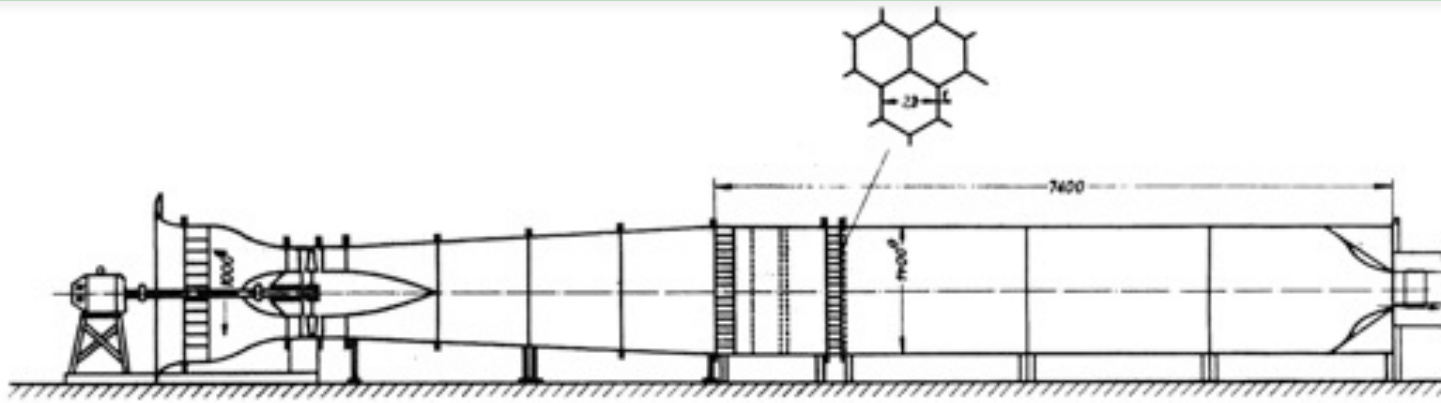


OUR GRID

- We (*uniquely*) have:
- independent paddles
 - feedback-control of angle



M = 163 mm

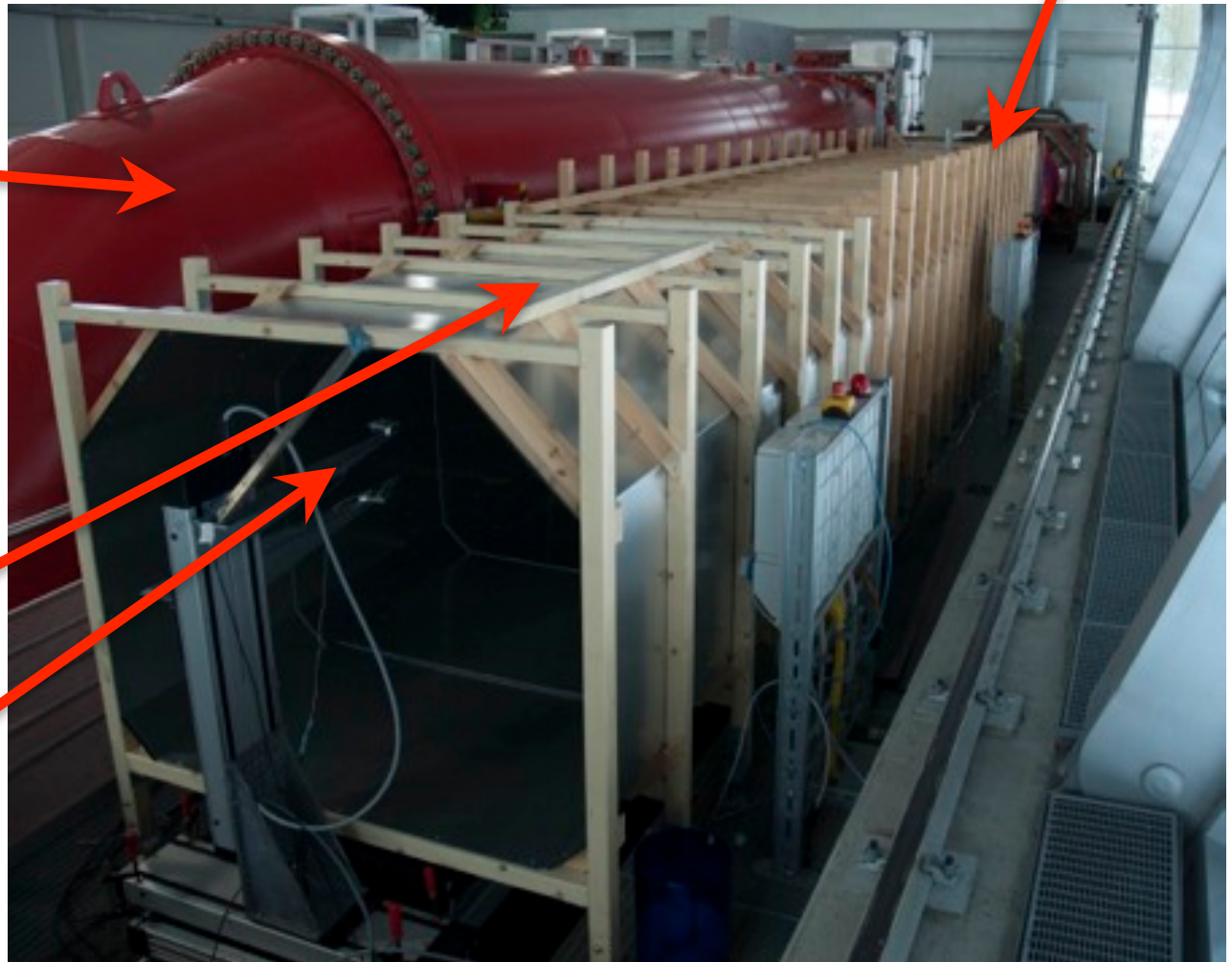


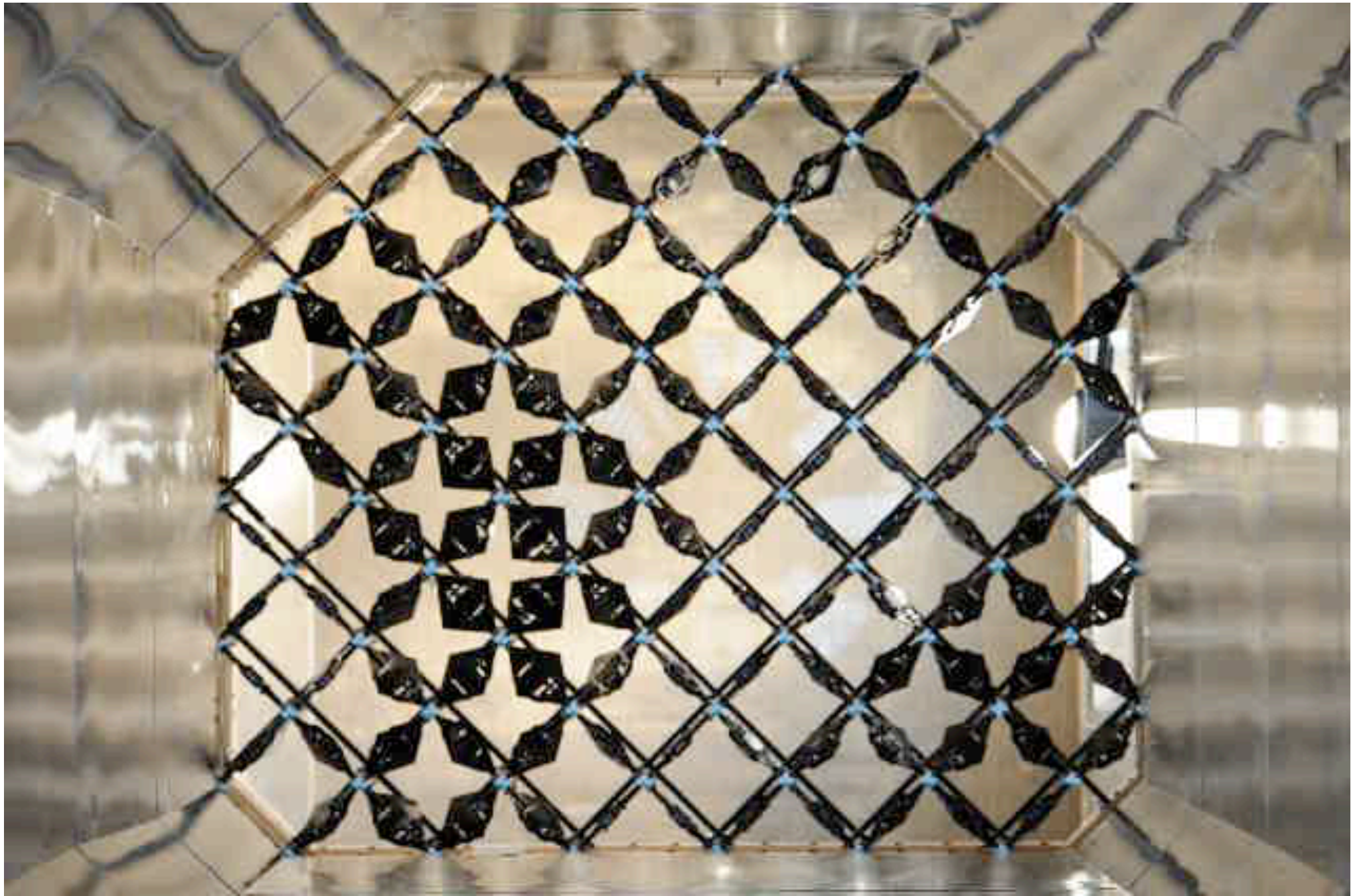
THE ACTIVE GRID

THE VDTT

THE PRANDTL TUNNEL
(1938)

HOT WIRE PROBES
50M DOWNSTREAM







Thank you:

E. Bodenschatz

Active grid

E. Cekli

F. Köhler

J. Kassel

F. Lachaussée

H. Grajewski

Wind tunnel

M. Sinhuber

H. Eckelmann

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E.-W. Saw

P.-Y. Lim

G. Good

D. Ivanov

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T. Schneider

J. Vollmer

H. Xu

A. Kopp

A. Kubitzek

O. Kurre

A. Renner

U. Schminke *et al.*